

Triangular solution to the general relativistic three-body problem

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Contents

- The background
- Equilateral triangular solution in GR
- **Triangular solution** in GR: general masses
- Summary

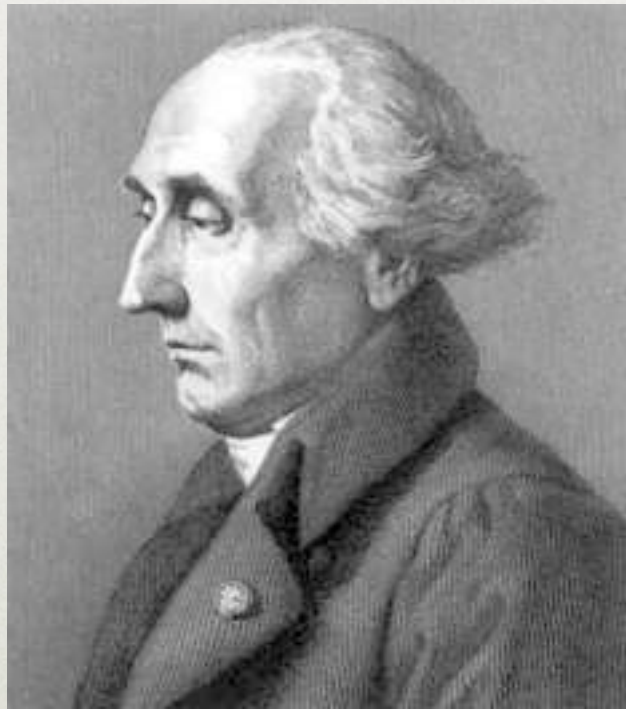
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Three-body problem

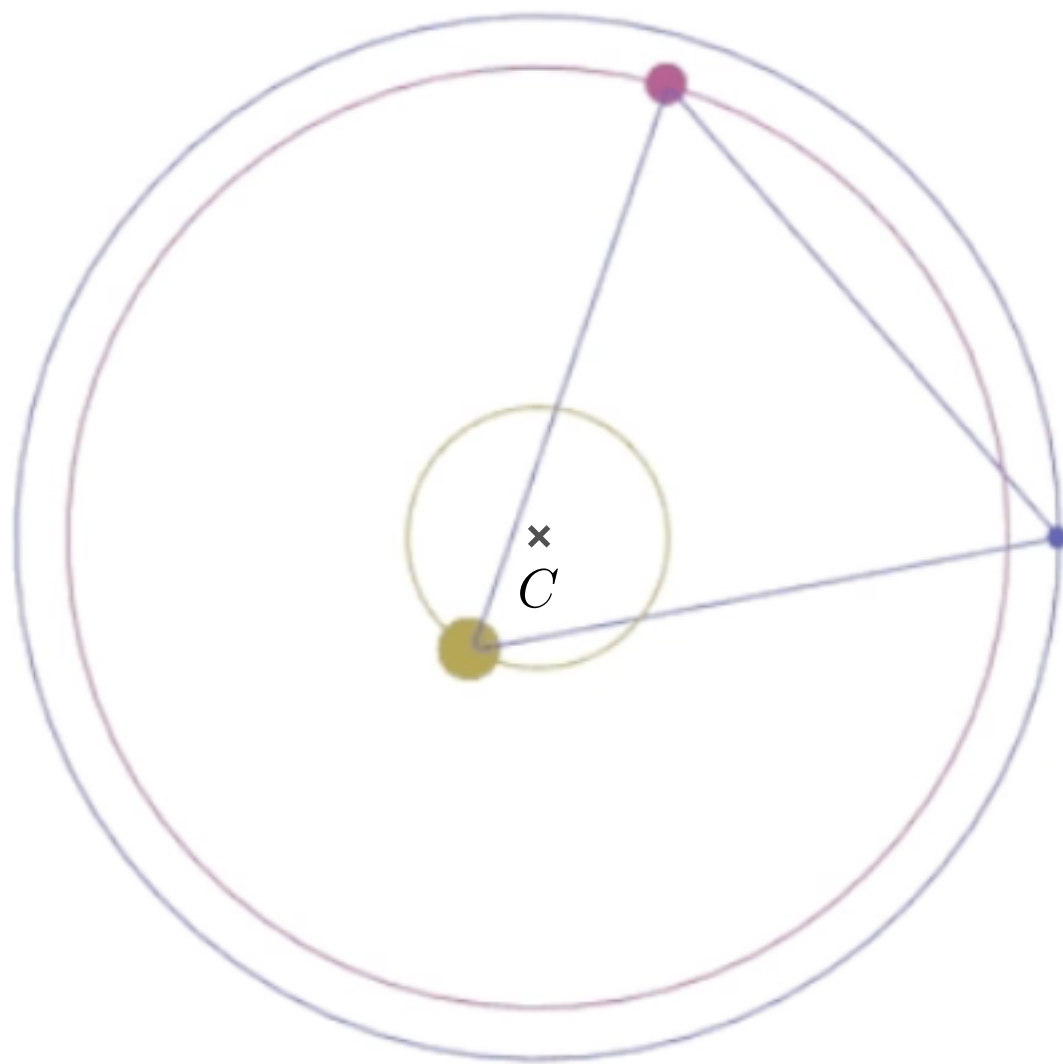
Particular solutions to the three-body problem

Euler's **collinear solution** (1765)
&



Lagrange's **equilateral triangular solution** (1772)

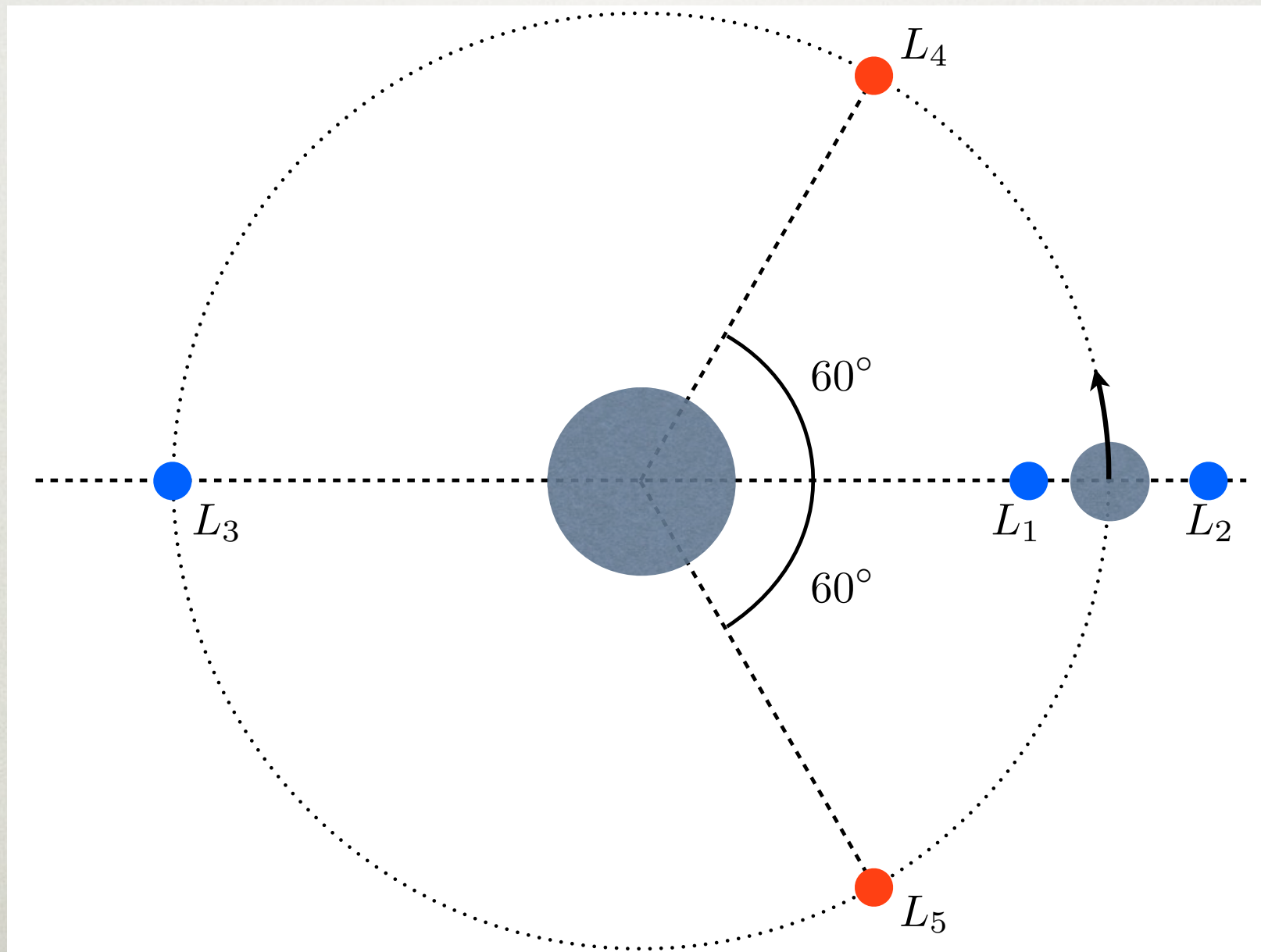
Equilateral triangular solution



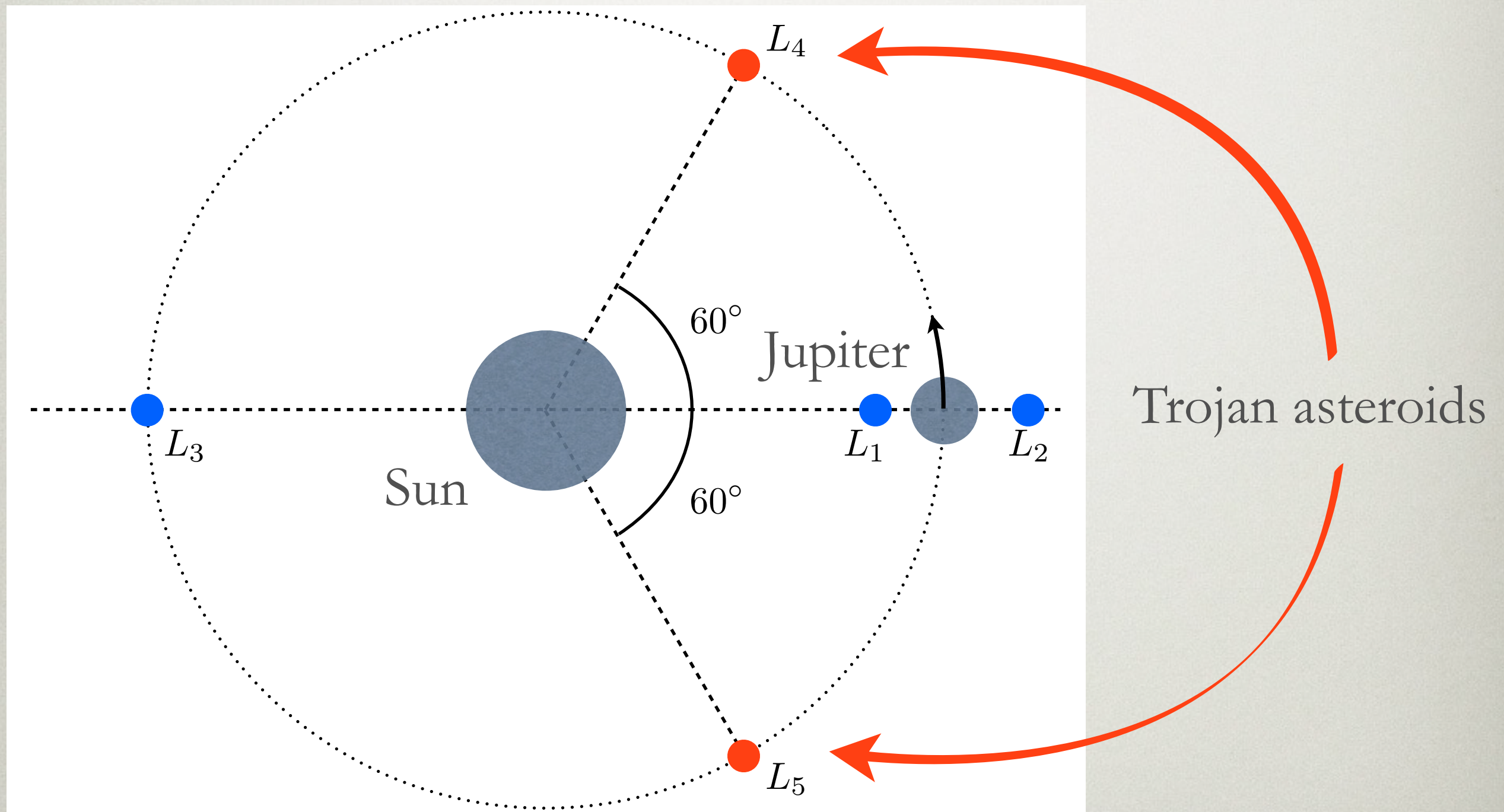
C : the center of mass

Lagrange points

L4 & L5



Lagrange points

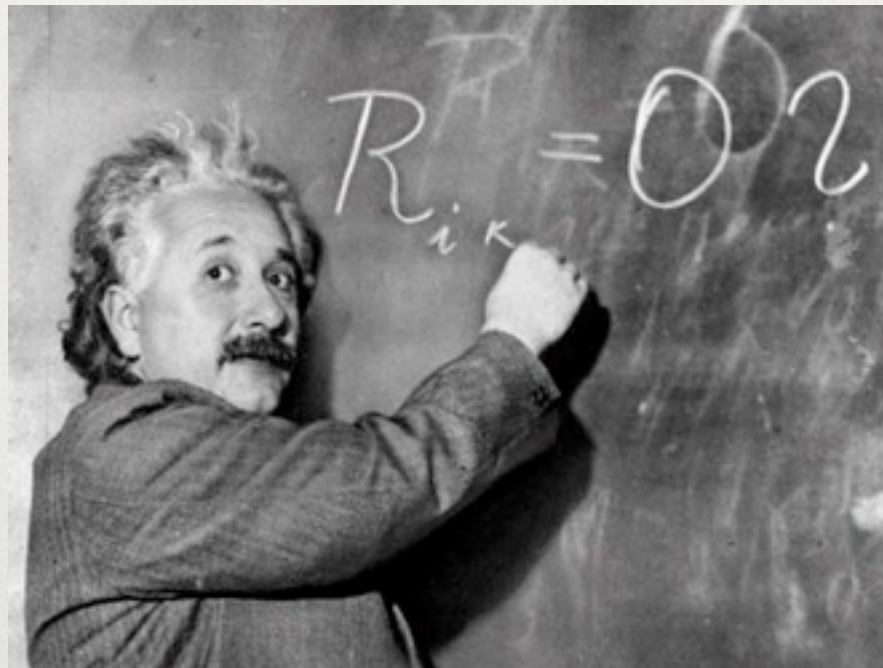


What happens in the general relativity (GR)?

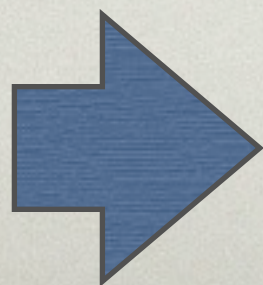
GR effects of Solar system

Two-body system: ○

(e.g. the perihelion precession of Mercury)



Three-body system: ?



It is interesting as a new test of GR

Gravity in GR

Einstein equation

$$G_{\mu\nu} = T_{\mu\nu}$$

↖ Matter

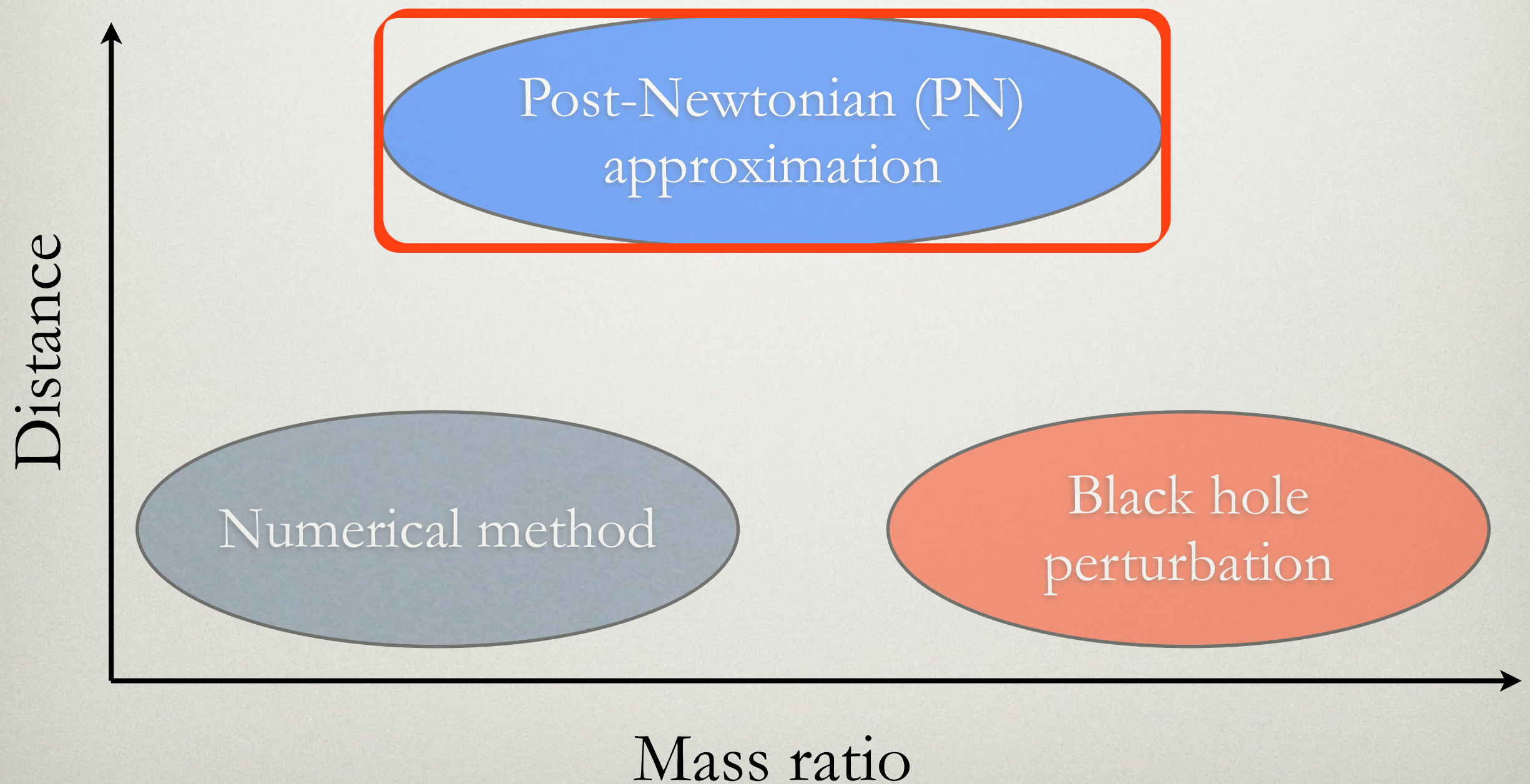
↗ Curvature of Spacetime



Gravity

10 non-linear second-order
partial differential equations

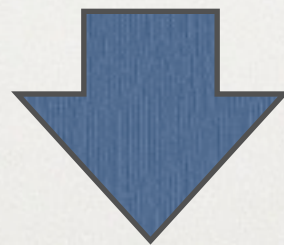
Post-Newtonian approximation



Post-Newtonian approximation

$$\lambda = \left(\frac{v}{c}\right)^2 \sim \frac{GM}{rc^2} \ll 1$$

v : velocity, c : light speed, r : distance, M : mass



Expand Einstein equation

Full order of GR

$$\underline{Q} = \underbrace{Q_0}_{\text{Newtonian term}} + \underbrace{Q_1\lambda + Q_2\lambda^2 + \dots}_{\text{corrections of GR}}$$

Newtonian term

First order of $\lambda = 1\text{PN}$ approximation

EIH equation of motion

Einstein-Infeld-Hoffman (EIH) equation of motion for N bodies

$$\begin{aligned}
 m_K \frac{d^2 \mathbf{r}_K}{dt^2} = & \sum_{A \neq K} \mathbf{r}_{AK} \frac{Gm_A m_K}{r_{AK}^3} \left[\underbrace{1}_{\text{Newtonian term}} - 4 \sum_{B \neq K} \underbrace{\frac{Gm_B}{c^2 r_{BK}}}_{\text{GR correction by mass}} - \sum_{C \neq A} \frac{Gm_C}{c^2 r_{CA}} \left(1 - \frac{\mathbf{r}_{AK} \cdot \mathbf{r}_{CA}}{2r_{CA}^2} \right) \right. \\
 & \left. + \underbrace{\left(\frac{\mathbf{v}_K}{c} \right)^2}_{\text{GR correction by velocity}} + 2 \left(\frac{\mathbf{v}_A}{c} \right)^2 - 4 \left(\frac{\mathbf{v}_A}{c} \right) \cdot \left(\frac{\mathbf{v}_K}{c} \right) - \frac{3}{2} \left(\frac{\left(\frac{\mathbf{v}_A}{c} \right) \cdot \mathbf{r}_{AK}}{r_{AK}} \right)^2 \right. \\
 & \left. - \sum_{A \neq K} \left[\left(\frac{\mathbf{v}_A}{c} \right) - \left(\frac{\mathbf{v}_K}{c} \right) \right] \frac{Gm_A m_K}{r_{AK}^3} \mathbf{r}_{AK} \cdot \left[3 \left(\frac{\mathbf{v}_A}{c} \right) - 4 \left(\frac{\mathbf{v}_K}{c} \right) \right] \right. \\
 & \left. + \frac{7}{2} \sum_{A \neq K} \sum_{C \neq A} \mathbf{r}_{CA} \underbrace{\frac{Gm_C m_K}{r_{CA}^3} \frac{Gm_A}{c^2 r_{AK}}}_{\text{Triple product}} \right]
 \end{aligned}$$

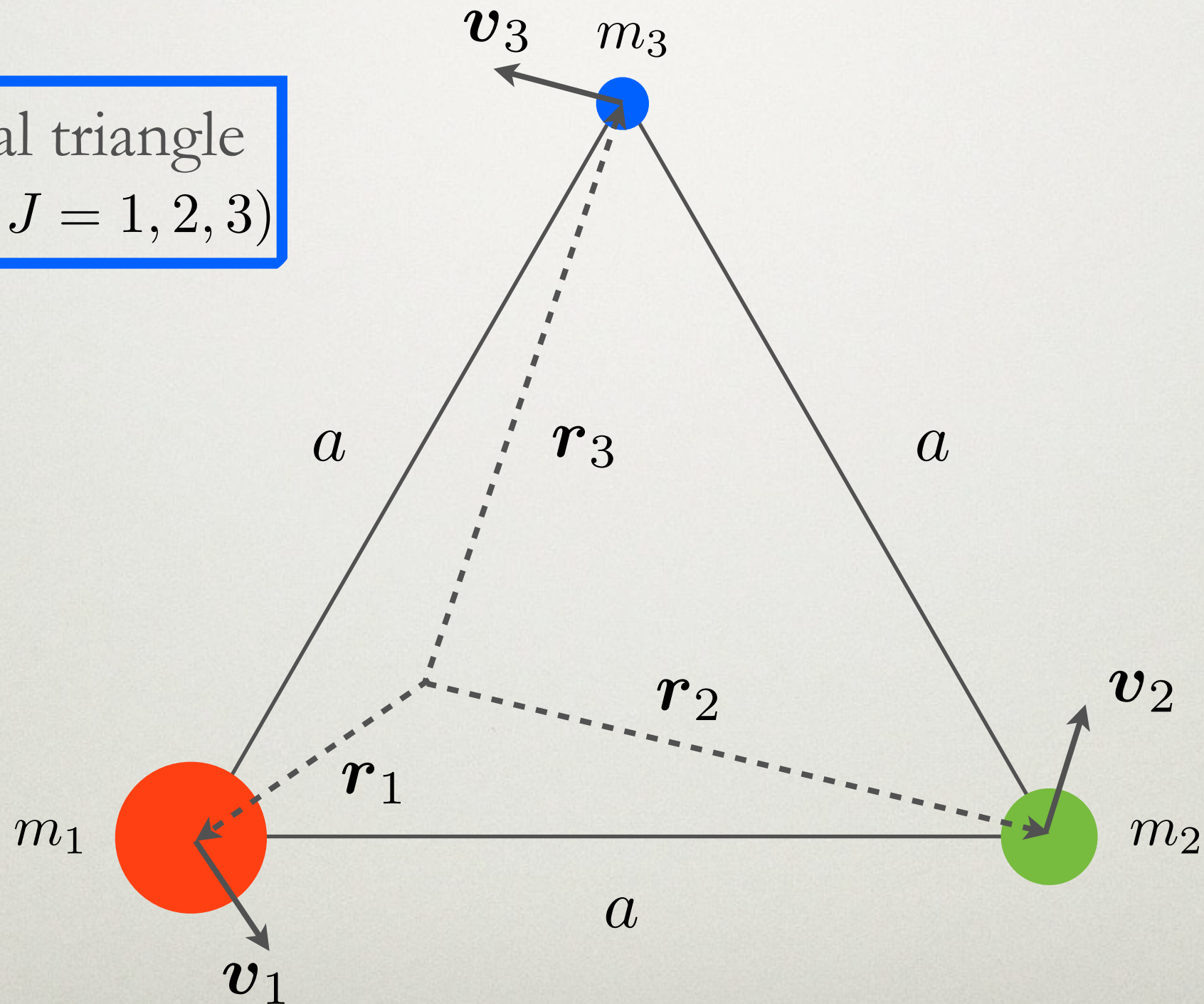
We look for an equilibrium solution in a circular motion

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Equilateral triangular configuration

Equilateral triangle
 $r_{IJ} = a \ (I, J = 1, 2, 3)$

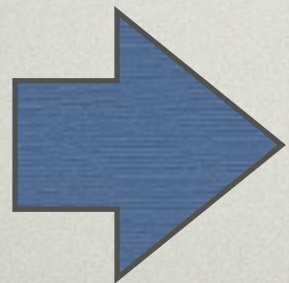


Center of mass at 1PN

$$\mathbf{r}_G = \frac{\sum_A \nu_A \mathbf{r}_A \left[1 + \frac{1}{2} \left(\left(\frac{v_A}{c} \right)^2 - \sum_{B \neq A} \frac{Gm_B}{c^2 r_{AB}} \right) \lambda \right]}{\sum_C \nu_C \left[1 + \frac{1}{2} \left(\left(\frac{v_C}{c} \right)^2 - \sum_{D \neq C} \frac{Gm_D}{c^2 r_{CD}} \right) \lambda \right]}$$

In general, it is different from Newtonian one

But



In this case, it agrees with the Newtonian one

Equilateral triangular solution at the 1PN

At 1PN order, EOM for m_1 becomes

$$-\omega^2 \mathbf{n}_1 = -\frac{M}{a^3} \mathbf{n}_1 + g_{PN1} \mathbf{n}_1 + \frac{\sqrt{3}}{16} \frac{M}{a^3} \frac{\nu_2 \nu_3 (\nu_2 - \nu_3)}{\nu_2^2 + \nu_2 \nu_3 + \nu_3^2} [6 + 9(\nu_2 + \nu_3)] \lambda \mathbf{n}_{\perp 1}$$

$$g_{PN1} = \frac{\lambda}{16(\nu_2^2 + \nu_2 \nu_3 + \nu_3^2)} \frac{M}{a^3} \times \left[48(\nu_2^2 + \nu_2 \nu_3 + \nu_3^2) - 2(8\nu_2^3 + 7\nu_2^2 \nu_3 + 7\nu_2 \nu_3^2 + 8\nu_3^3) + (16\nu_2^4 + 41\nu_2^3 \nu_3 + 84\nu_2^2 \nu_3^2 + 41\nu_2 \nu_3^3 + 16\nu_3^4) \right]$$

$$\omega : \text{angular velocity, } \nu_I \equiv m_I/M, \quad M = \sum_I m_I \quad (I = 1, 2, 3)$$

$$\mathbf{n}_1 \equiv \mathbf{r}_1/|\mathbf{r}_1|, \quad \mathbf{n}_{\perp 1} \equiv \mathbf{v}_1/|\mathbf{v}_1|, \quad \mathbf{n}_{\perp 1} \text{ is normal to } \mathbf{n}_1$$

Equilateral triangular solution at the 1PN

$$-\omega^2 \mathbf{n}_1 = -\frac{M}{a^3} \mathbf{n}_1 + g_{PN1} \mathbf{n}_1 + \frac{\sqrt{3}}{16} \frac{M}{a^3} \frac{\nu_2 \nu_3 (\nu_2 - \nu_3)}{\nu_2^2 + \nu_2 \nu_3 + \nu_3^2} [6 + 9(\nu_2 + \nu_3)] \lambda \mathbf{n}_{\perp 1}$$

In only **2** cases, bodies satisfy EOM;

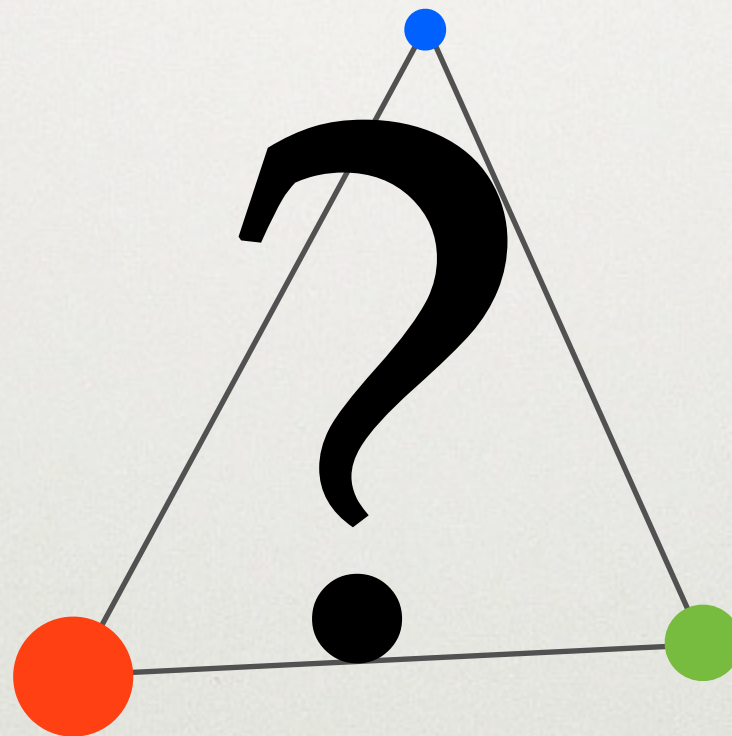
- mass ratio 1 : 1 : 1
- mass ratio 0 : 0 : 1

This solution does not always exist in GR

[Ichita, KY & Asada, PRD **83**, 084026 (2011)]

Equilateral triangular solution at the 1PN

For **the arbitrary mass ratio**,
the solution exist?



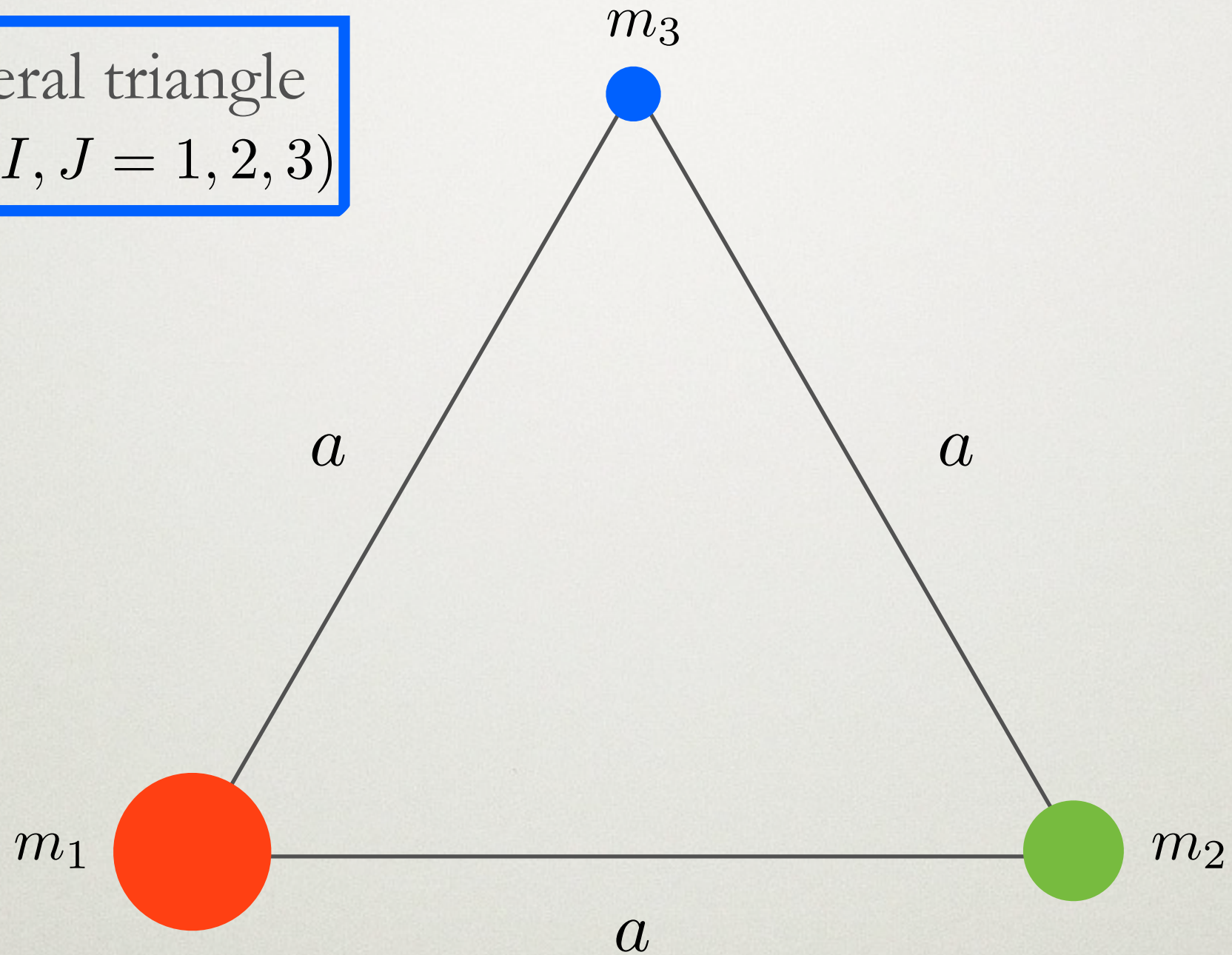
cf. [E. Krefetz, Astron. J. **72**, 471 (1967)]
for **restricted** 3-body problem

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Corrections of distance

Equilateral triangle
 $r_{IJ} = a \ (I, J = 1, 2, 3)$

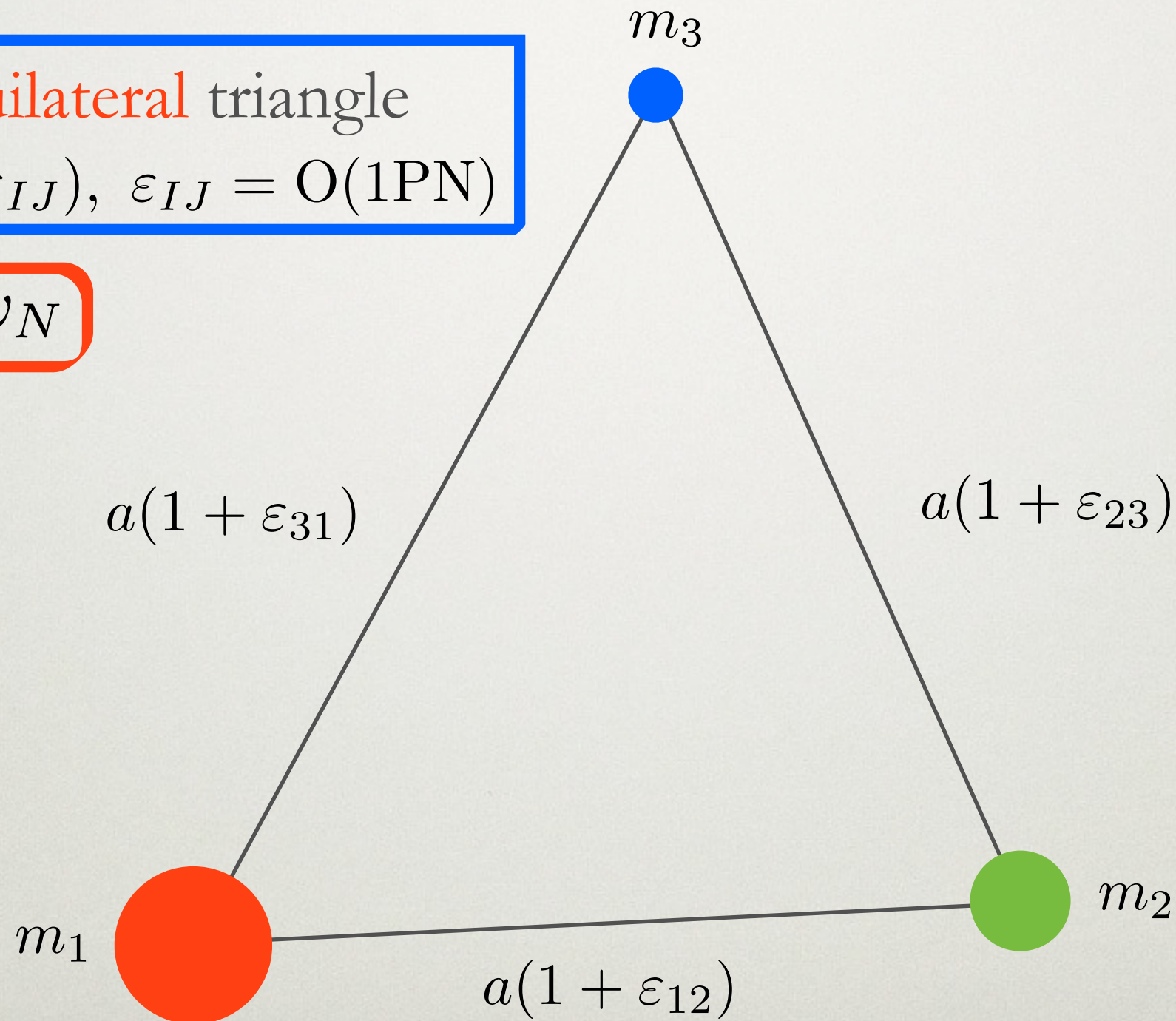


Corrections of distance

PN **inequilateral** triangle

$$r_{IJ} = a(1 + \varepsilon_{IJ}), \quad \varepsilon_{IJ} = \mathcal{O}(1\text{PN})$$

$$\omega = \omega_N$$



We can ignore the correction for center of mass

Triangular solution at the 1PN

EOM for m_1 becomes

$$\begin{aligned} -\omega^2 \mathbf{r}_1 = & -\omega_N^2 \mathbf{r}_1 \\ & + \nu_2 \left(-3 + \nu_1 \nu_2 + \nu_2 \nu_3 + \nu_3 \nu_1 - \frac{3}{8} \nu_3 [5 - 3(\nu_1 + \nu_2)] \right) \lambda \mathbf{r}_{21} \\ & + \nu_3 \left(-3 + \nu_1 \nu_2 + \nu_2 \nu_3 + \nu_3 \nu_1 - \frac{3}{8} \nu_2 [5 - 3(\nu_3 + \nu_1)] \right) \lambda \mathbf{r}_{31} \\ & - 3(\nu_2 \varepsilon_{12} \mathbf{r}_{21} + \nu_3 \varepsilon_{31} \mathbf{r}_{31}) \end{aligned}$$

$$\boxed{\omega = \omega_N} \quad \Rightarrow \quad \boxed{} = 0$$

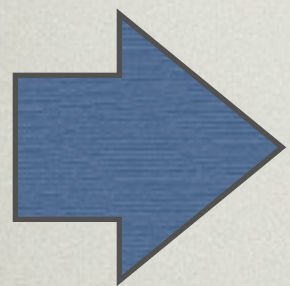
Triangular solution at the 1PN

As a result, we could uniquely express ε_{IJ}

$$\varepsilon_{12} = - \left[1 - \frac{1}{3}(\nu_1\nu_2 + \nu_2\nu_3 + \nu_3\nu_1) + \frac{1}{8}\nu_3[5 - 3(\nu_1 + \nu_2)] \right] \lambda,$$

$$\varepsilon_{23} = - \left[1 - \frac{1}{3}(\nu_1\nu_2 + \nu_2\nu_3 + \nu_3\nu_1) + \frac{1}{8}\nu_1[5 - 3(\nu_2 + \nu_3)] \right] \lambda,$$

$$\varepsilon_{31} = - \left[1 - \frac{1}{3}(\nu_1\nu_2 + \nu_2\nu_3 + \nu_3\nu_1) + \frac{1}{8}\nu_2[5 - 3(\nu_1 + \nu_3)] \right] \lambda.$$



Triangular solution for the arbitrary mass ratio at 1PN

[KY, H. Asada, in prep.]

Application for Solar system

Corrections for L4 (L5) of Solar system [m]

Planet	Sun-Planet	Sun-L4 (L5)	Planet-L4 (L5)
Earth	-1477	-1477	-1477 -923
Jupiter	-1477	-1477	-1477 -922

The sign + denotes increase of distance

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Summary

- We found a **triangular solution** with 1PN corrections of distance
- The 1PN triangle is smaller than the Newtonian one (for same masses)
- This result may be applied to also near SMBH and compact binary
- The future observations are needed

Future works

- The Stability
- The Gravitational wave
- An elliptical motion
- More bodies
- Higher order PN approximation

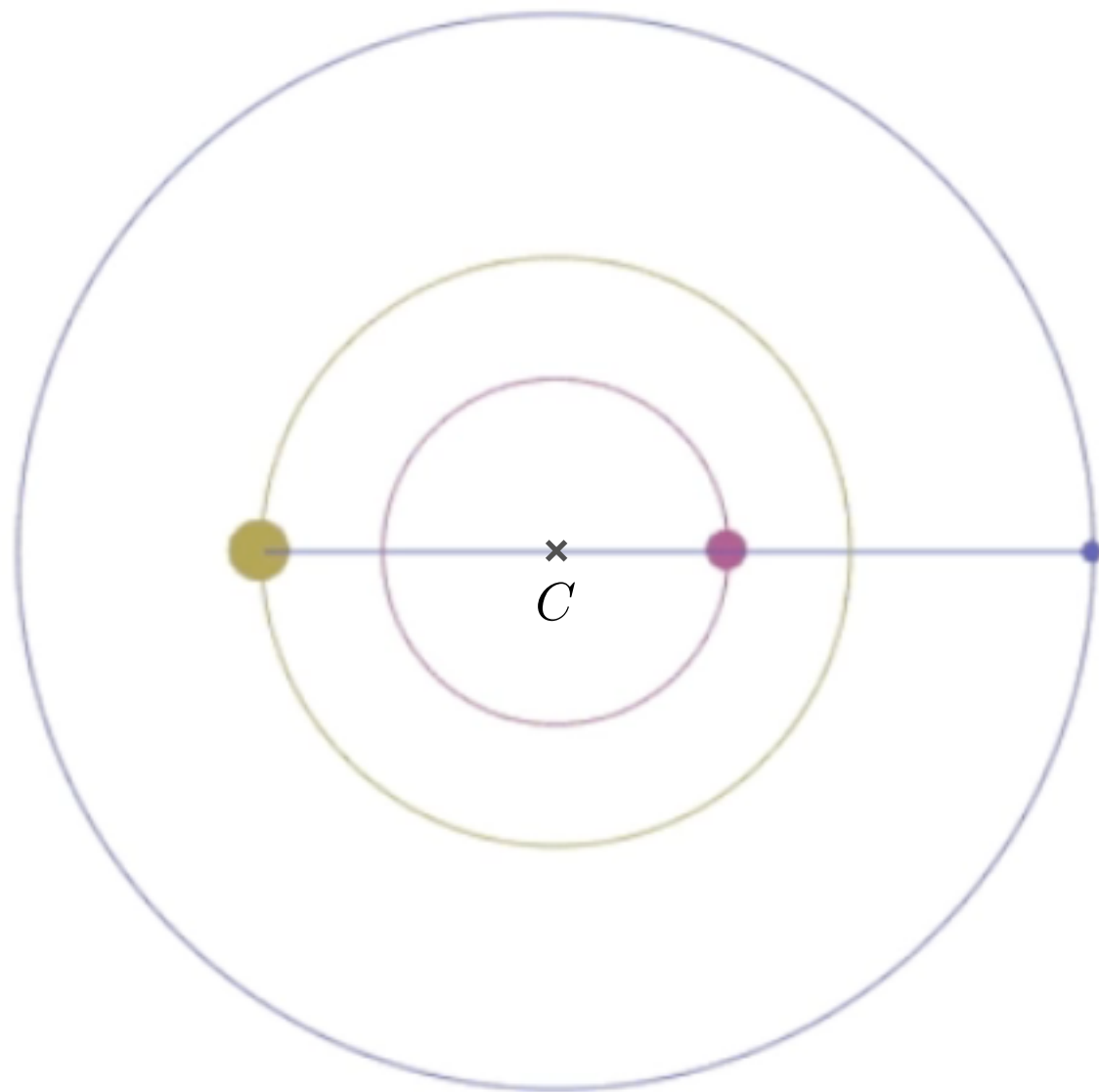


Thank you for your attention

Appendix

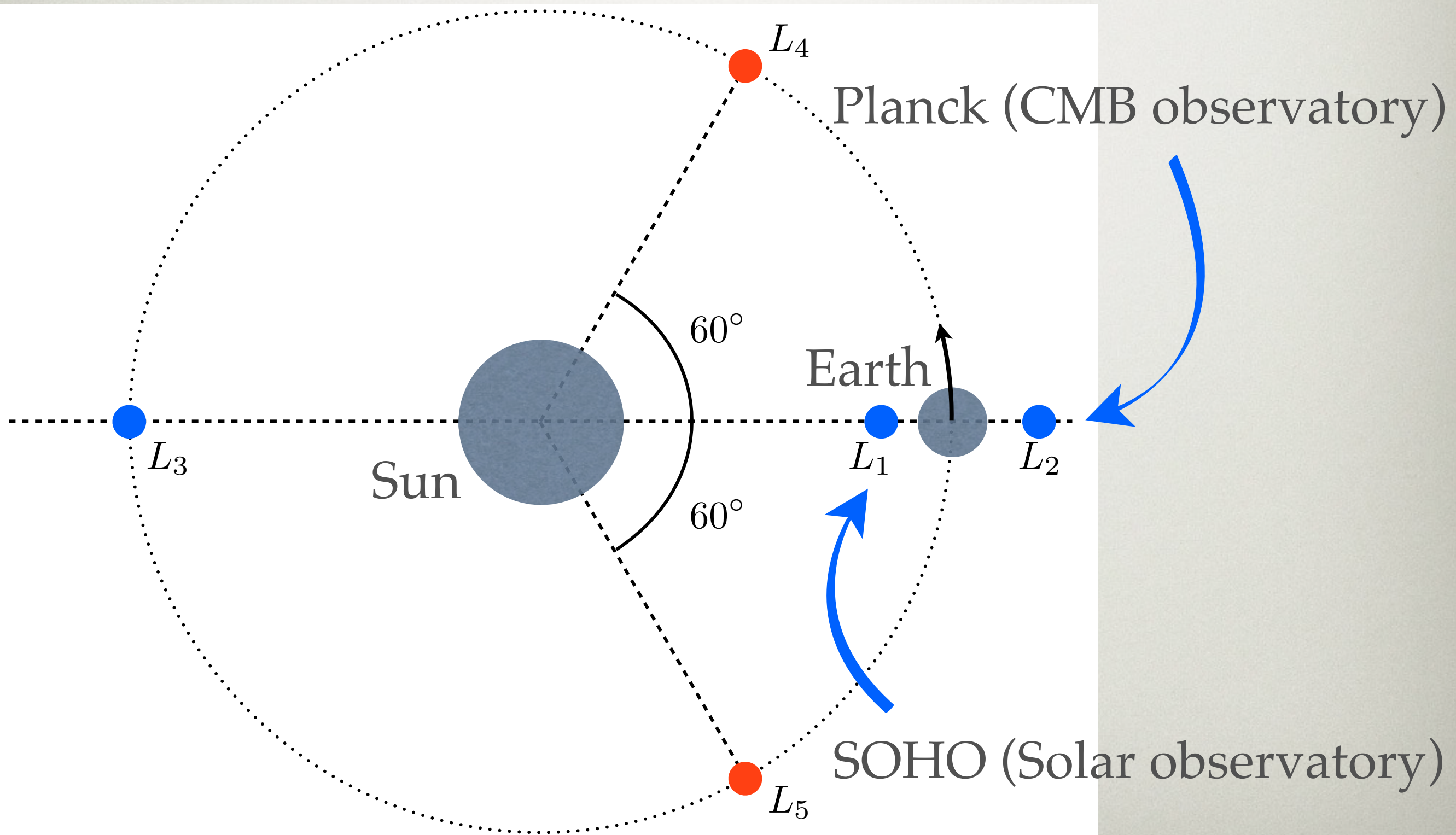
Collinear solution

Collinear solution



C : 共通重心

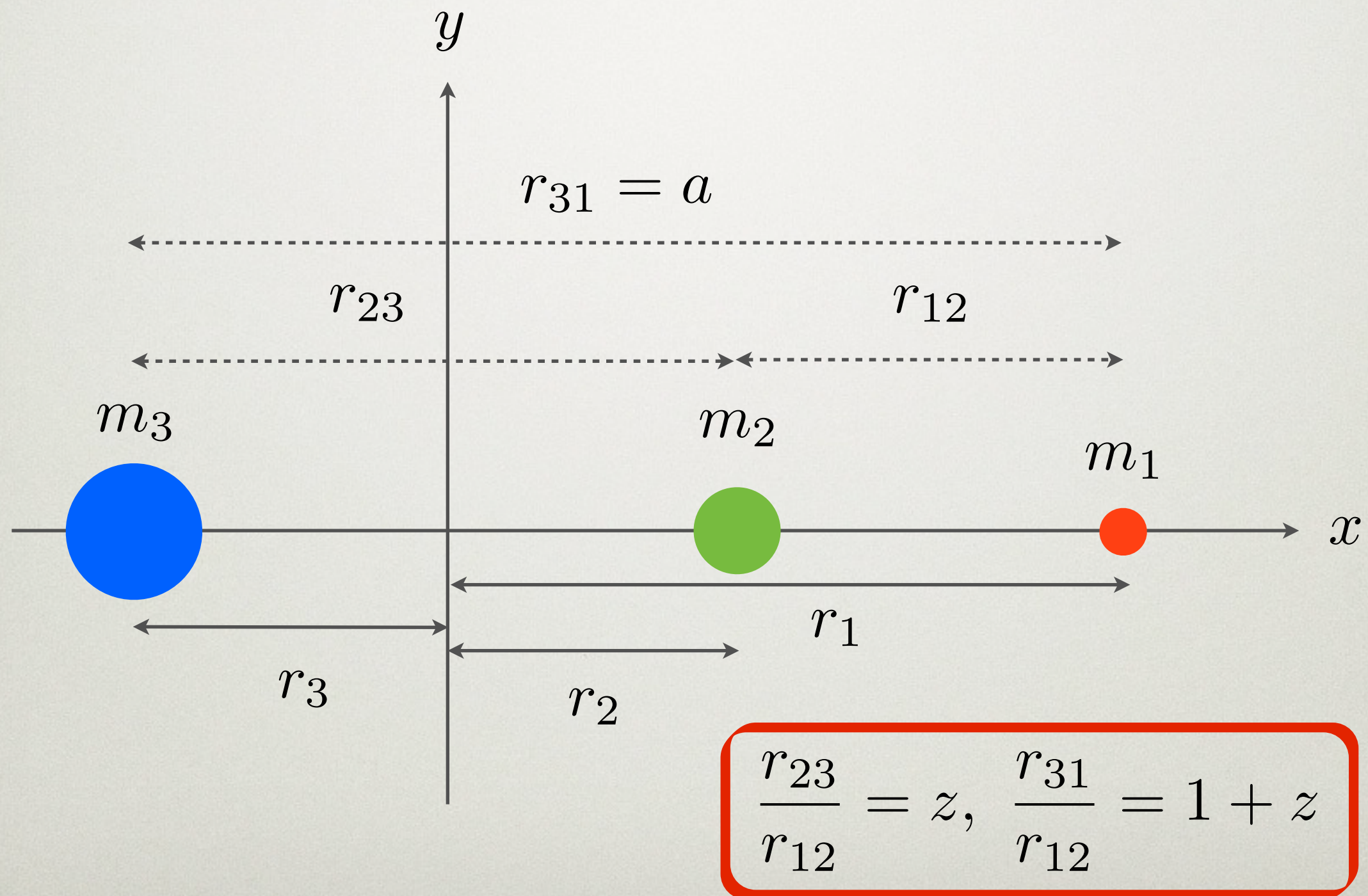
Lagrange points



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- Three-body problem in Newtonian gravity
- **Triangular solution** to the general relativistic three-body problem
- **Collinear solution** to the general relativistic three-body problem
- Application to Solar system
- Summary

Collinear configuration



Collinear solution in Newtonian

$$\frac{r_{23}}{r_{12}} = z, \quad \frac{r_{31}}{r_{12}} = 1 + z$$

Equations of motion for 3 masses

$$r_1 \omega^2 = \left[\frac{m_2}{r_{12}^2} + \frac{m_3}{r_{31}^2} \right], \quad r_2 \omega^2 = \left[-\frac{m_1}{r_{12}^2} + \frac{m_3}{r_{23}^2} \right], \quad r_3 \omega^2 = \left[-\frac{m_1}{r_{31}^2} - \frac{m_2}{r_{23}^2} \right]$$



$$(m_1 + m_2)z^5 + (3m_1 + 2m_2)z^4 + (3m_1 + m_2)z^3 - (m_2 + 3m_3)z^2 - (2m_2 + 3m_3)z - (m_2 + m_3) = 0$$

Collinear solution at 1PN

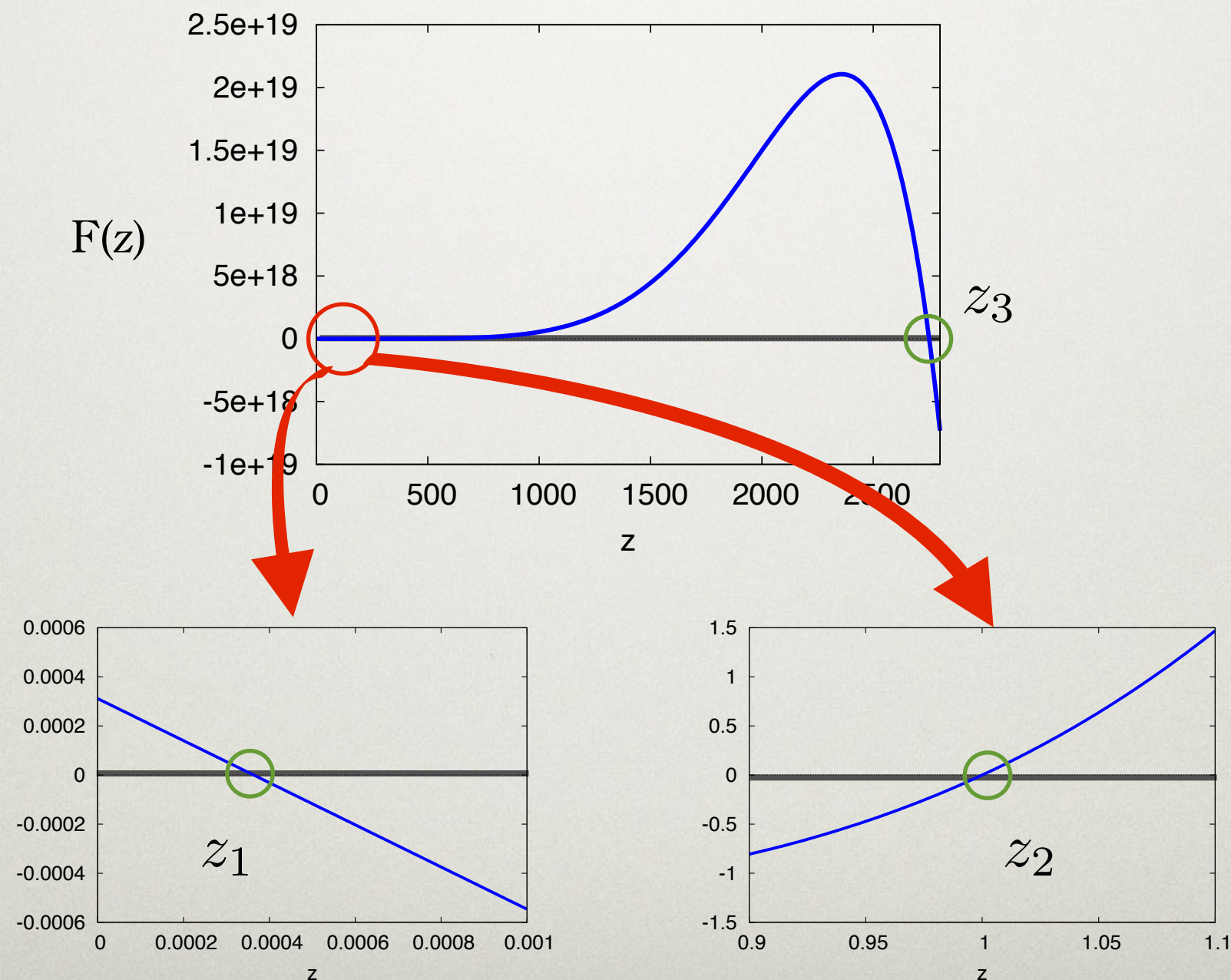
$$F(z) \equiv \sum_{k=0}^7 A_k z^k = 0$$

$$\begin{aligned} A_7 &= \frac{GM}{ac^2} \left[-4 - 2(\nu_1 - 4\nu_3) + 2(\nu_1^2 + 2\nu_1\nu_3 - 2\nu_3^2) - 2\nu_1\nu_3(\nu_1 + \nu_3) \right], & A_3 &= -(1 - \nu_1 + 2\nu_3) + \frac{GM}{ac^2} \left[6 + 2(2\nu_1 + 5\nu_3) - 4(4\nu_1^2 + \nu_1\nu_3 - 2\nu_3^2) \right. \\ & & & \left. + 2(3\nu_1^3 + 2\nu_1^2\nu_3 - \nu_1\nu_3^2 - 3\nu_3^3) \right], \\ A_6 &= 1 - \nu_3 + \frac{GM}{ac^2} \left[-13 - (10\nu_1 - 17\nu_3) + 2(2\nu_1^2 + 8\nu_1\nu_3 - \nu_3^2) \right. \\ & & \left. + 2(\nu_1^3 - 2\nu_1^2\nu_3 - 3\nu_1\nu_3^2 - \nu_3^3) \right], & A_2 &= -(2 - 2\nu_1 + \nu_3) + \frac{GM}{ac^2} \left[15 - (5\nu_1 - 18\nu_3) - 4(4\nu_1^2 + 5\nu_1\nu_3) \right. \\ & & & \left. + 6(\nu_1^3 + \nu_1^2\nu_3 - \nu_3^3) \right], \\ A_5 &= 2 + \nu_1 - 2\nu_3 + \frac{GM}{ac^2} \left[-15 - (18\nu_1 - 5\nu_3) + 4(5\nu_1\nu_3 + 4\nu_3^2) \right. \\ & & \left. + 6(\nu_1^3 - \nu_1\nu_3^2 - \nu_3^3) \right], & A_1 &= -(1 - \nu_1) + \frac{GM}{ac^2} \left[13 - (17\nu_1 - 10\nu_3) + 2(\nu_1^2 - 8\nu_1\nu_3 - 2\nu_3^2) \right. \\ & & & \left. + 2(\nu_1^3 + 3\nu_1^2\nu_3 + 2\nu_1\nu_3^2 - \nu_3^3) \right], \\ A_4 &= 1 + 2\nu_1 - \nu_3 + \frac{GM}{ac^2} \left[-6 - 2(5\nu_1 + 2\nu_3) - 4(2\nu_1^2 - \nu_1\nu_3 - 4\nu_3^2) \right. \\ & & \left. + 2(3\nu_1^3 + \nu_1^2\nu_3 - 2\nu_1\nu_3^2 - 3\nu_3^3) \right], & A_0 &= \frac{GM}{ac^2} \left[4 - 2(4\nu_1 - \nu_3) + 2(2\nu_1^2 - 2\nu_1\nu_3 - \nu_3^2) + 2\nu_1\nu_3(\nu_1 + \nu_3) \right]. \end{aligned}$$

[Yamada & Asada, PRD 82, 104019 (2010)]

7th-order equation

$$\nu_1 = 1/7, \nu_2 = 5/7, \nu_3 = 1/7, \lambda = GM/c^2 a = 10^4$$



7th-order equation

7th-order equation may have 3 positive roots



Smallest root and largest one give fast motion,
which is comparable to speed of light.



of physical roots = 1

PN angular velocity

PN angular velocity is expressed as

$$\omega = \omega_N \left(1 + \frac{F_M}{2F_N} + \frac{F_V}{2R_{31}} \right),$$

$$F_N = \frac{M}{a^2 z^2} \left[(\nu_1 + \nu_3) z^2 + (1 - \nu_1 - \nu_3)(1 + z^2)(1 + z)^2 \right],$$
$$F_M = -\frac{M^2}{a^3 z^3} \left[(4 - 4\nu_1 + \nu_3)(1 - \nu_1 - \nu_3) \right. \\ + (12 - 7\nu_1 + 3\nu_3)(1 - \nu_1 - \nu_3)z \\ + (12 - \nu_1 + \nu_3)(1 - \nu_1 - \nu_3)z^2 \\ + (8 - 7\nu_1 - 7\nu_3 + 8\nu_1\nu_3 + 3\nu_1^2 + 3\nu_3^2)z^3 \\ + (12 + \nu_1 - \nu_3)(1 - \nu_1 - \nu_3)z^4 \\ + (12 + 3\nu_1 - 7\nu_3)(1 - \nu_1 - \nu_3)z^5 \\ \left. + (4 + \nu_1 - 4\nu_3)(1 - \nu_1 - \nu_3)z^6 \right],$$
$$F_V = \frac{M}{(1 + z)^2 z^2} \left[-\nu_1^2(1 - \nu_1 - \nu_3) \right. \\ - 2\nu_1(1 + \nu_1 - \nu_3)(1 - \nu_1 - \nu_3)z \\ + (2 - 2\nu_1 + \nu_3 + 6\nu_1\nu_3 - 3\nu_3^2 + \nu_1^3 - 3\nu_1^2\nu_3 - 3\nu_1\nu_3^2 + \nu_3^3)z^2 \\ + 2(2 - \nu_1 - \nu_3)(1 + \nu_1 + \nu_3 - \nu_1^2 + \nu_1\nu_3 - \nu_3^2)z^3 \\ + (2 + \nu_1 - 2\nu_3 - 3\nu_1^2 + 6\nu_1\nu_3 + \nu_1^3 - 3\nu_1^2\nu_3 - 3\nu_1\nu_3^2 + \nu_3^3)z^4 \\ - 2\nu_3(1 - \nu_1 + \nu_3)(1 - \nu_1 - \nu_3)z^5 \\ \left. - \nu_3^2(1 - \nu_1 - \nu_3)z^6 \right].$$

ω_N : Newtonian angular velocity

PN angular velocity

For same masses and full length $a(= r_{31})$
we can prove inequality

$$\omega < \omega_N$$

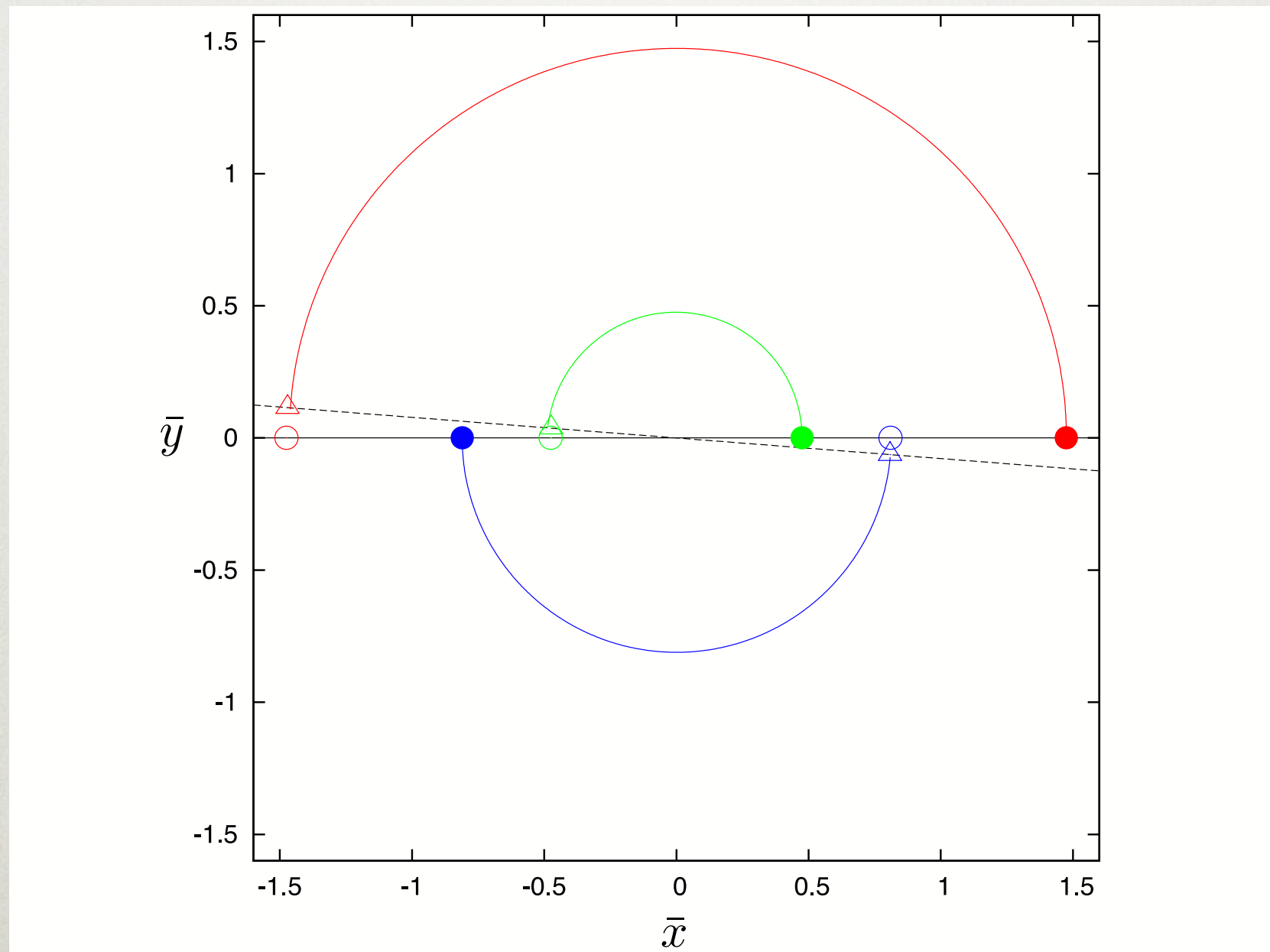
Thus, PN angular velocity is always smaller than Newtonian



Namely, PN orbital period is always longer than Newtonian

Numerical example

$$\nu_1 : \nu_2 : \nu_3 = 1 : 2 : 3, \quad \lambda = 10^{-2} \left(\frac{v}{c} \sim 10^{-1} \right)$$



[Yamada & Asada, PRD 83, 024040 (2011)]