

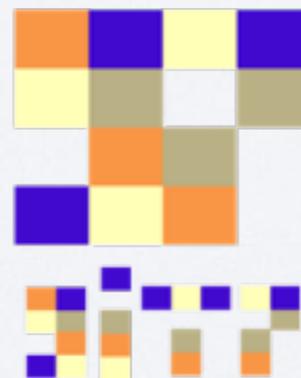
中心エンジン駆動超新星エジェクタの 多次元シミュレーション

Multi-dimensional simulations of
central-engine powered supernova ejecta

Akihiro Suzuki (YITP)

collaborator: Keiichi Maeda (Kyoto U.)

Suzuki & Maeda (2017) MNRAS 466 2633 and recent updates



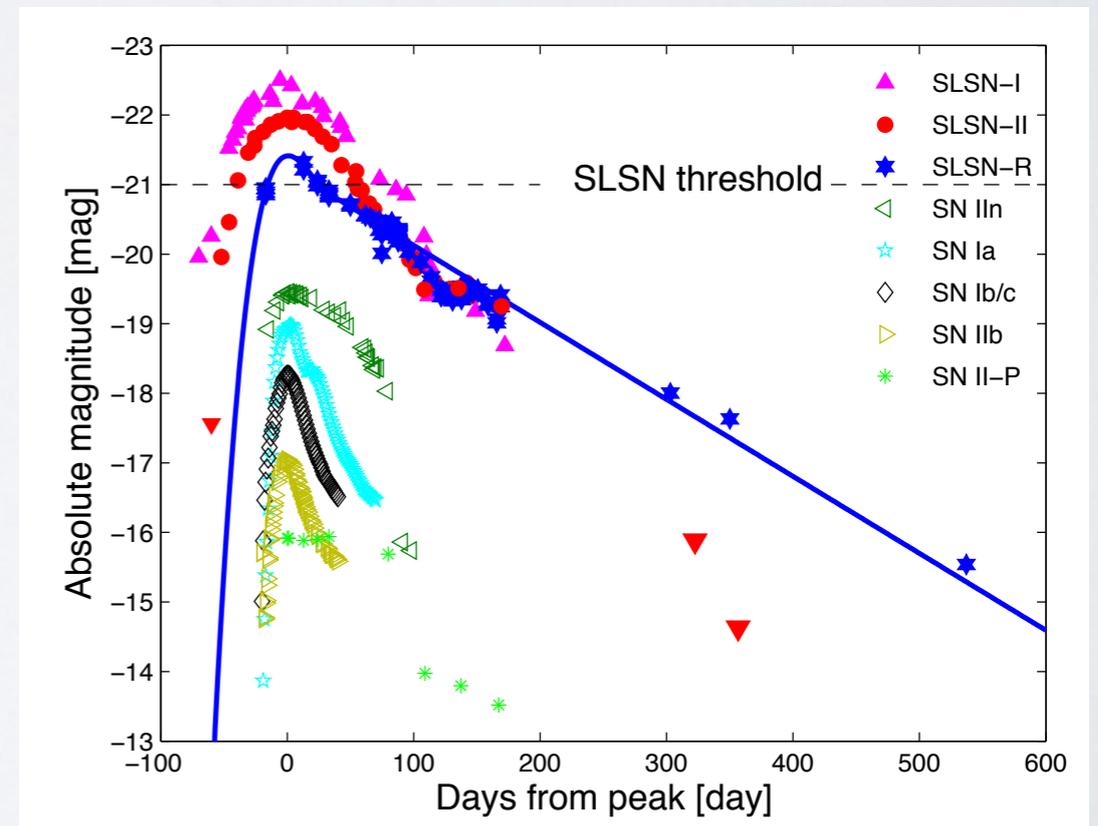
Ordinary and Extra-ordinary CCSNe

- CCSNe energetics: Canonically,
 - gravitational energy $E_{\text{grav}} \sim GM_{\text{ns}}^2/R_{\text{ns}} \sim 10^{53}$ [erg]
 - kinetic energy $E_{\text{kin}} \sim 1\%$ of $E_{\text{grav}} \sim 10^{51}$ [erg]
 - total radiated energy $E_{\text{rad}} \sim$ **less than 1%** of $E_{\text{kin}} \sim <10^{49}$ [erg]
 - ejecta mass: **a few - 10 M_{\odot}**
 - photospheric velocity: typically, $\sim 10,000$ [km/s]

- However, some unusual SNe have been found:
 - **broad-lined Ic SNe (Ic-BL)**: photospheric velocity larger by a factor of 2-3 \sim **a few 10^4 [km/s]**, which implies $E_{\text{kin}} \sim 10^{52}$ [erg] $> 10^{51}$ [erg]
 - **Superluminous SNe (SLSNe)**: $E_{\text{rad}} \sim 10^{51}$ [erg] $> 10^{49}$ [erg]

Superluminous SNe

- Superluminous supernovae (SLSNe): SNe **10-100 times brighter** than normal SNe (Quimby+2007, Barbary+2009 etc, see Gal-Yam+2012 for review)
- They are found by recent “unbiased” transient survey projects (e.g., Palomar transient factory, Pan-STARRS).
- The following classification based on their optical spectra has been proposed (analogy to standard SNe).
 - 1) **SLSN-I** : no Hydrogen feature
 - 2) **SLSN-II** : Hydrogen feature
 - 3) **SLSN-R** : subclass of SLSN-I, their light curves can be explained by the decay of radioactive ^{56}Ni (e.g., $3M_{\odot}$ Ni for SN 2007bi)
- Total radiated energy can be $\sim 10^{51}$ [erg] (~ explosion energy of normal CCSNe)

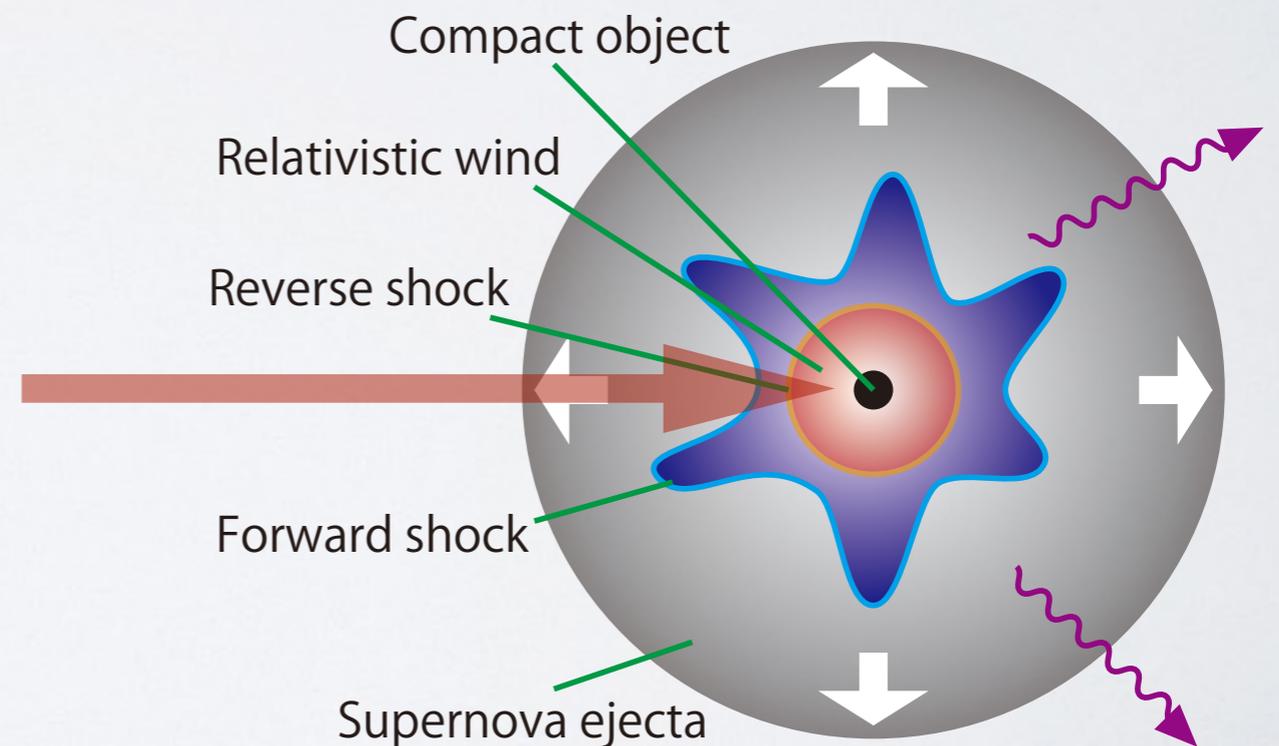
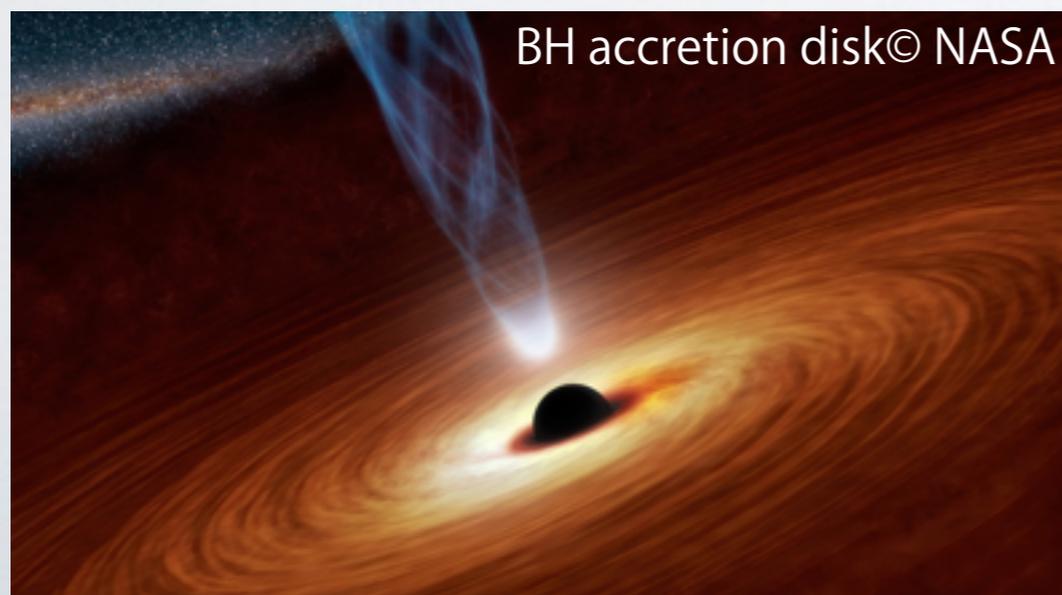


What is the origin of SLSNe-I?

↑ light curves of standard SNe, SLSNe (Gal-Yam 2012)

Proposed models and progenitors for SLSNe

- CSM interaction
- pair-instability SNe (very massive progenitor with $\sim 100\text{-}300M_{\odot}$ at ZAMS)
- **additional energy injection from the central engine** : magnetar spin-down (e.g., Kasen&Bildsten 2010, Woosley 2010) or BH accretion (Dexter&Kasen 2013)



mili-second Magnetar scenario

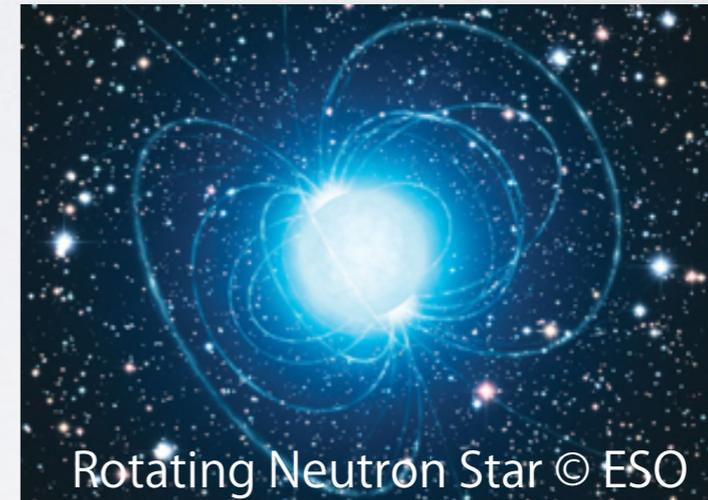
- After the gravitational collapse of the iron core, a massive star experience the core bounce and its outer layer with mass M_{ej} is expelled by neutrino-driven explosion with $E_{kin}=10^{51}$ [erg] (standard scenario for CCSNe).
- a neutron star with a strong dipole magnetic field is assumed to form immediately after the neutrino-driven explosion.

radius $R_{ns} \sim 10\text{km}$

moment of inertia $I_{ns} \sim 10^{45} \text{ g cm}^2$

initial period of $P_i \sim 1 \text{ ms}$

$E_{rot} = I_{ns}\Omega_i^2/2 \simeq 2 \times 10^{52} \text{ erg.}$



- spin-down of the new-born magnetar is expected to power the SN ejecta

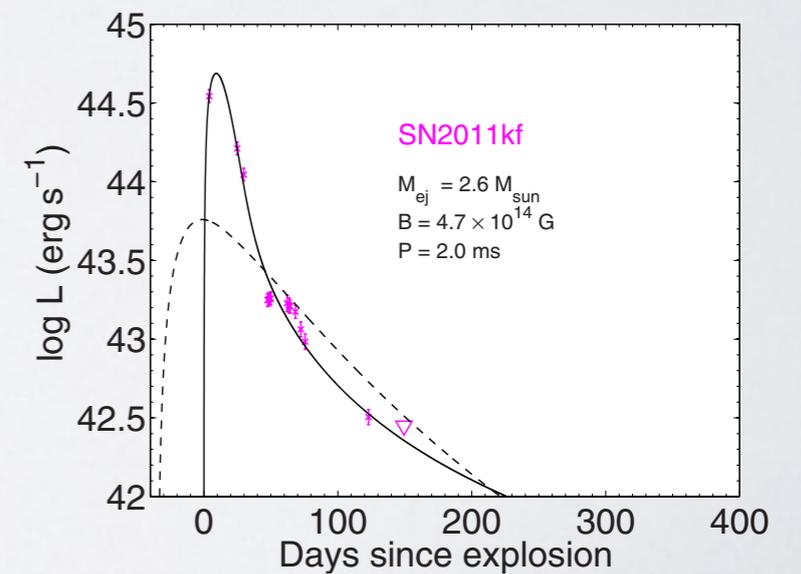
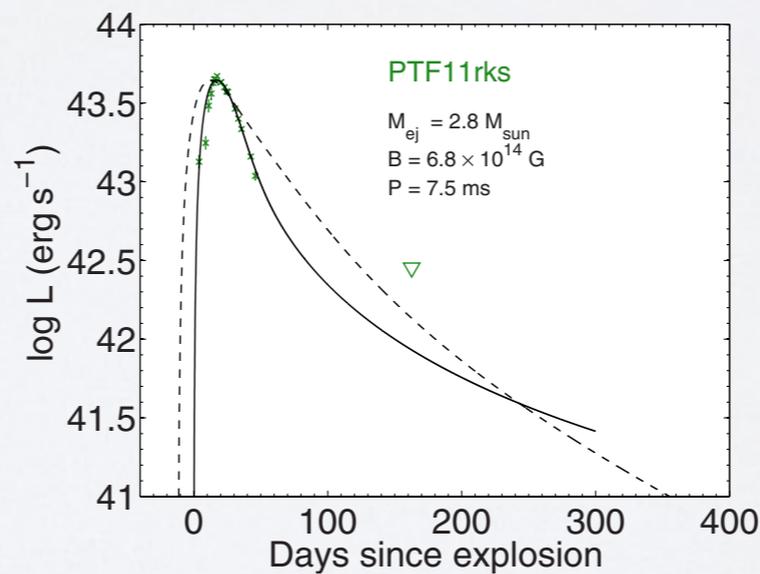
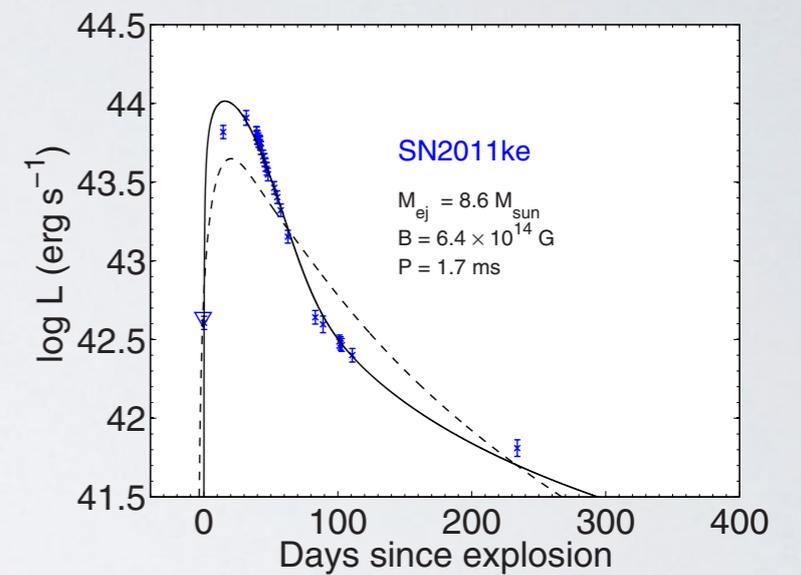
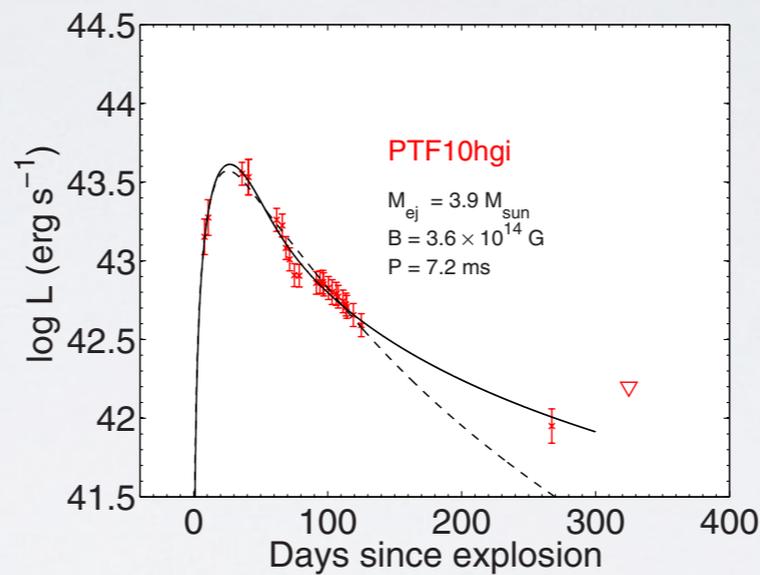
$$L = \frac{E_{rot}/t_{ch}}{(1 + t/t_{ch})^2}$$

$$L \simeq \frac{B^2 R_{ns}^6 \Omega_i^4}{6c^3} \sim 10^{49} B_{15}^2 R_{ns,6}^6 P_{i,-3}^{-4} \text{ erg s}^{-1}$$

$$t_{ch} = \frac{6I_{ns}c^3}{B^2 R_{ns}^6 \Omega_i^2} = 4.1 \times 10^3 I_{ns,45} B_{15}^2 R_{ns,6}^6 P_{i,-3}^2 \text{ s.}$$

mili-second Magnetar scenario

- one-box light curve model for SNe with magnetar energy injection
- LCs are explained by “tuning” several free parameters, M_{ej} , B , and P_i .
- Magnetar scenario looks successful when one-box model is considered.



↑ Magnetar model fit to SLSNe-I (Inserra+2013)

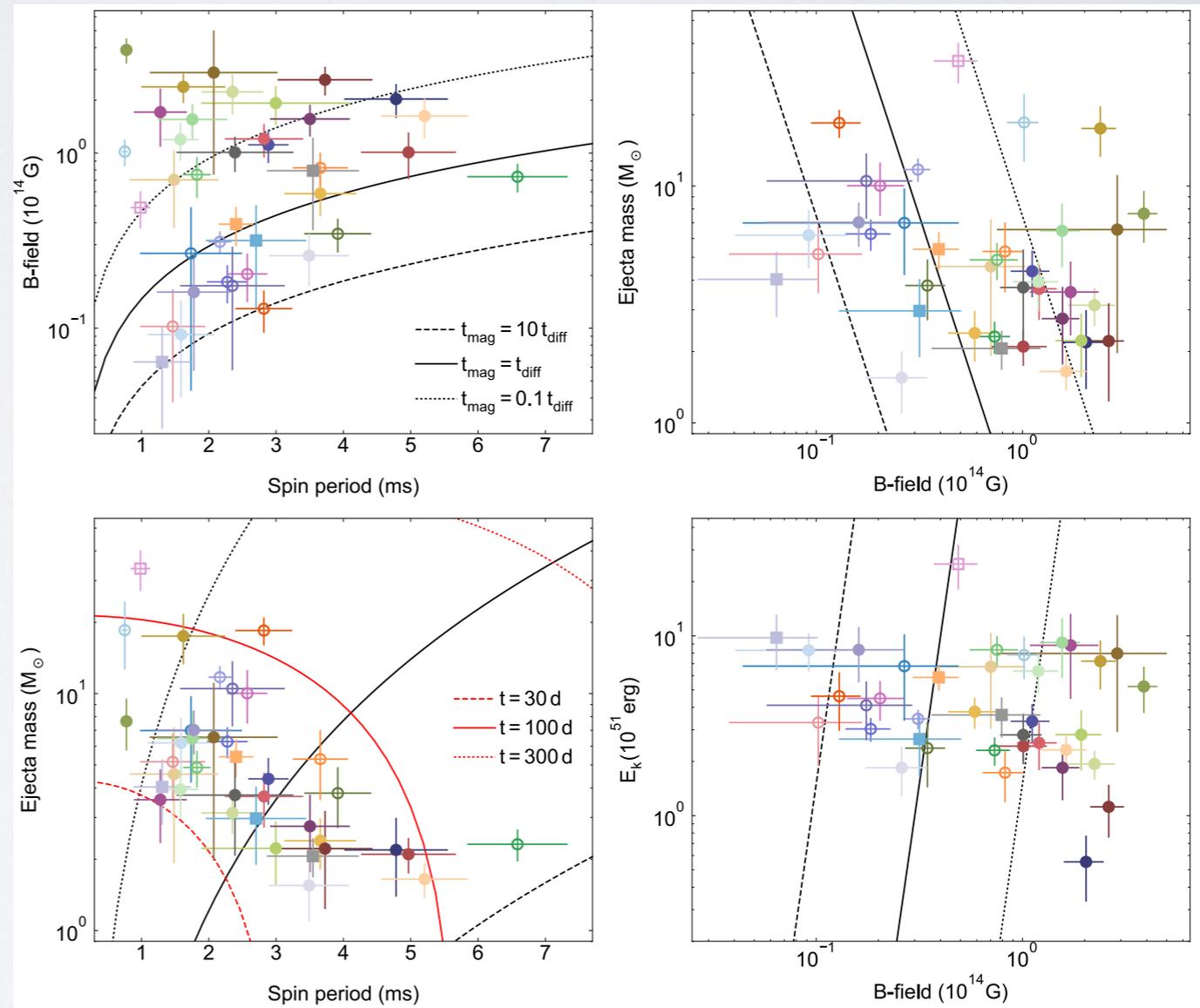
Magnetar scenario

- one-box light curve model for SNe with magnetar energy injection
- LCs are explained by “tuning” several free parameters, M_{ej} , B , and P_i .
- Magnetar scenario looks successful when one-box model is considered.

- Magnetar fit :

- spin-period $\sim 1 - 7$ [ms]
- $B \sim 10^{13} - \text{a few } 10^{14}$ [G]
- time-scale $\sim \text{a few } 10\text{-}100$ days
- $E_k \sim 10^{51} - 10^{52}$ [erg]
- $M_{ej} \sim 2 - 10 M_{\odot}$

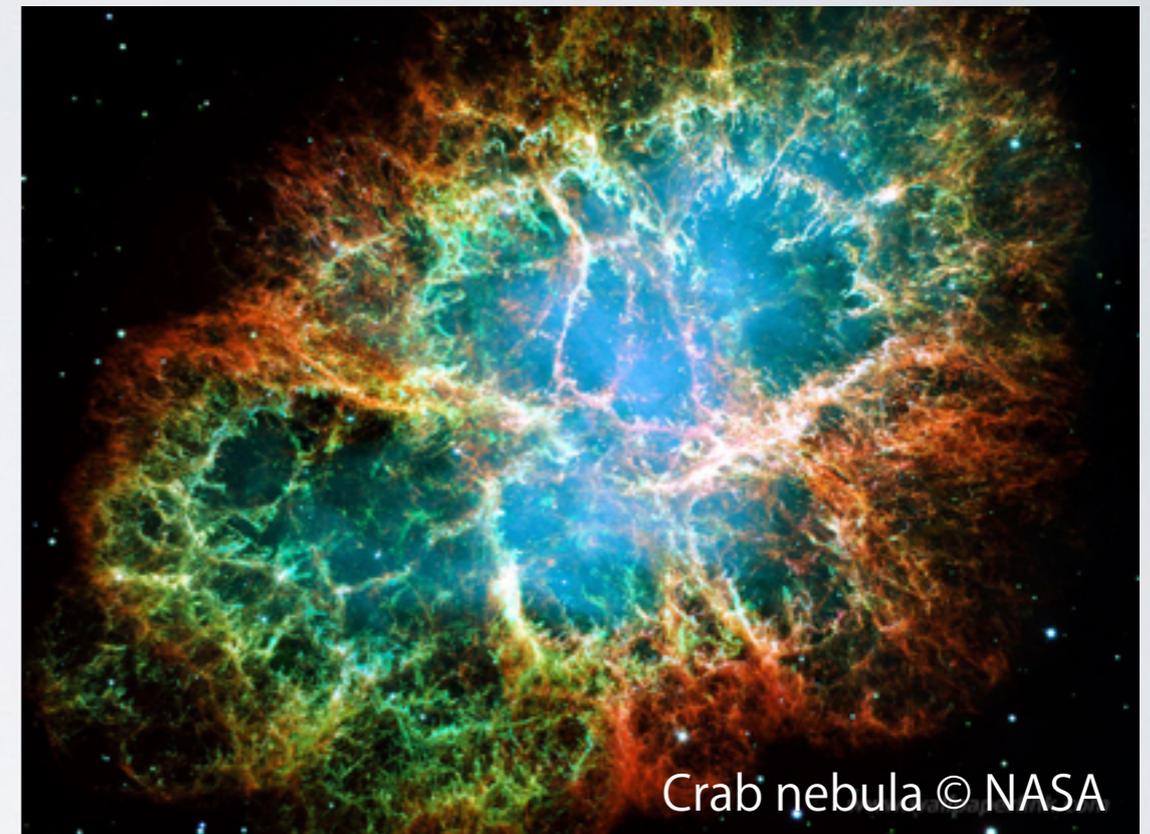
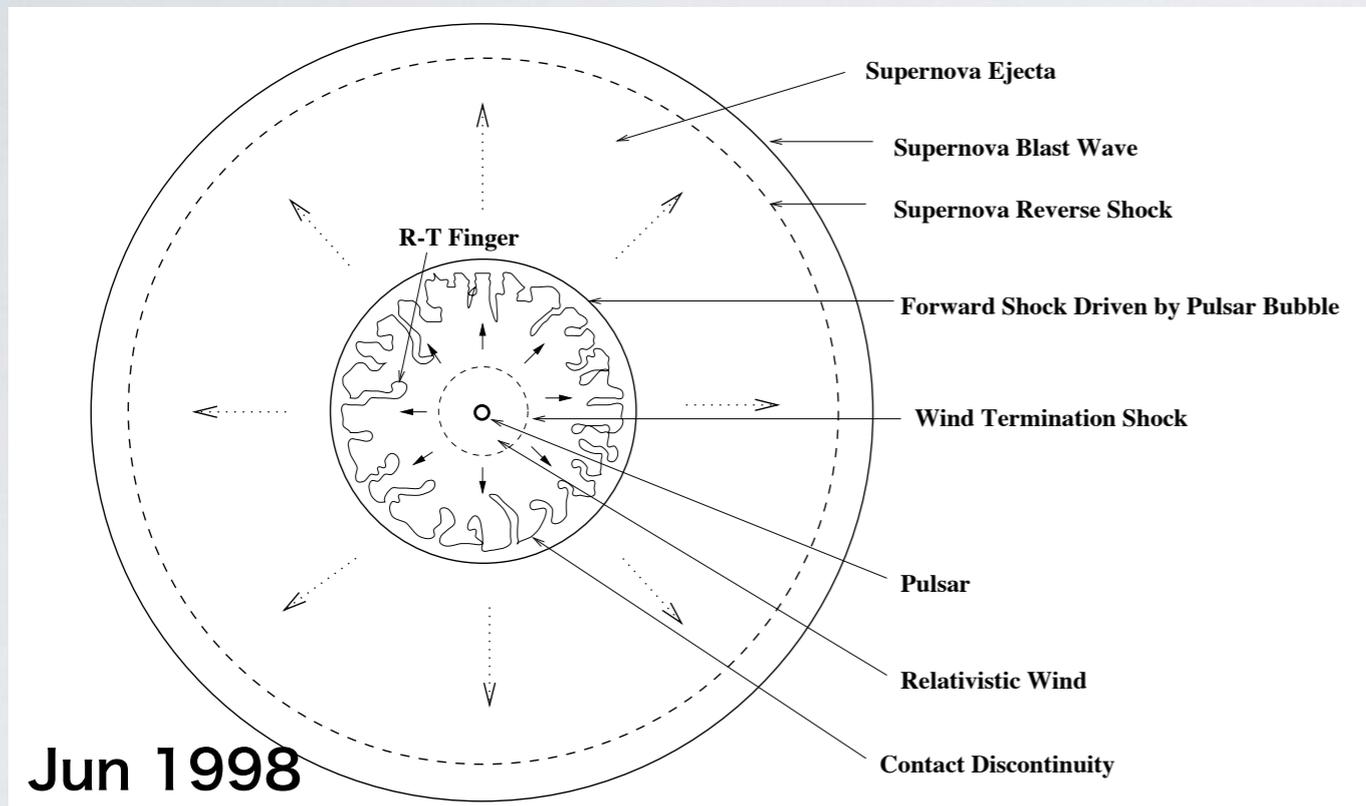
(e.g., Nicholl+2017)



↑ Magnetar model fit to SLSNe-I (Nicholl+2017)

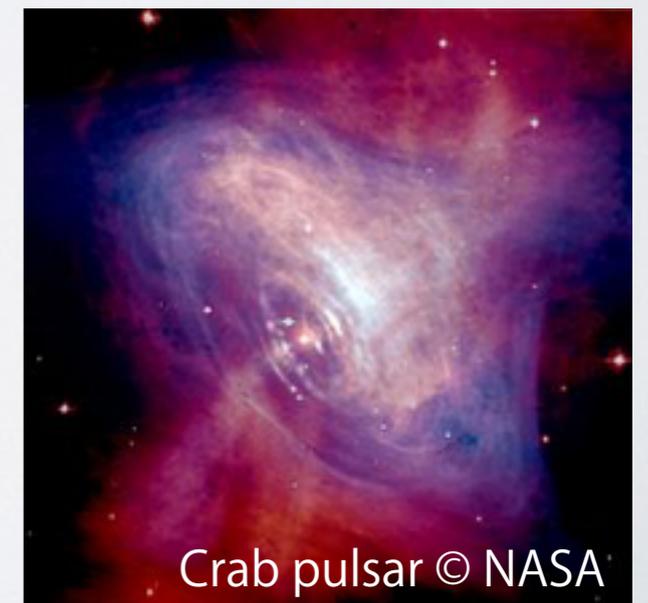
Relativistic wind from magnetized NS

- SN ejecta (or SNR) pushed by a pulsar wind nebula
- e.g., Crab nebula



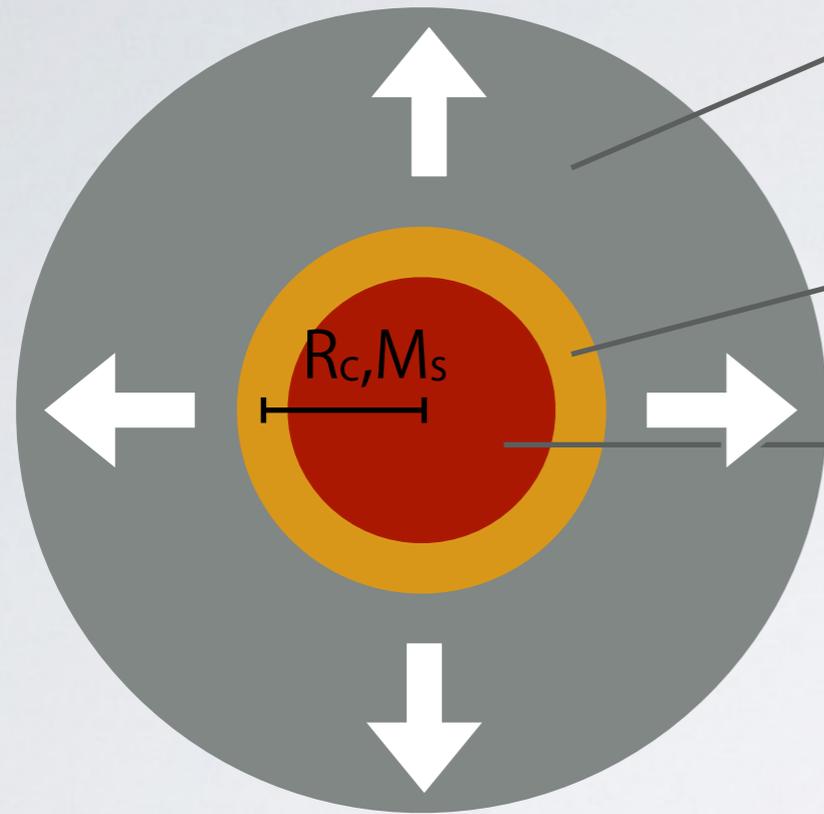
- galactic PWNe: injected energy $E_{inj} < SN$ explosion energy E_{exp}

What happens when $E_{inj} > E_{exp}$?
(or maybe $E_{inj} \gg E_{exp}$)



1D spherical picture of SN-wind interaction

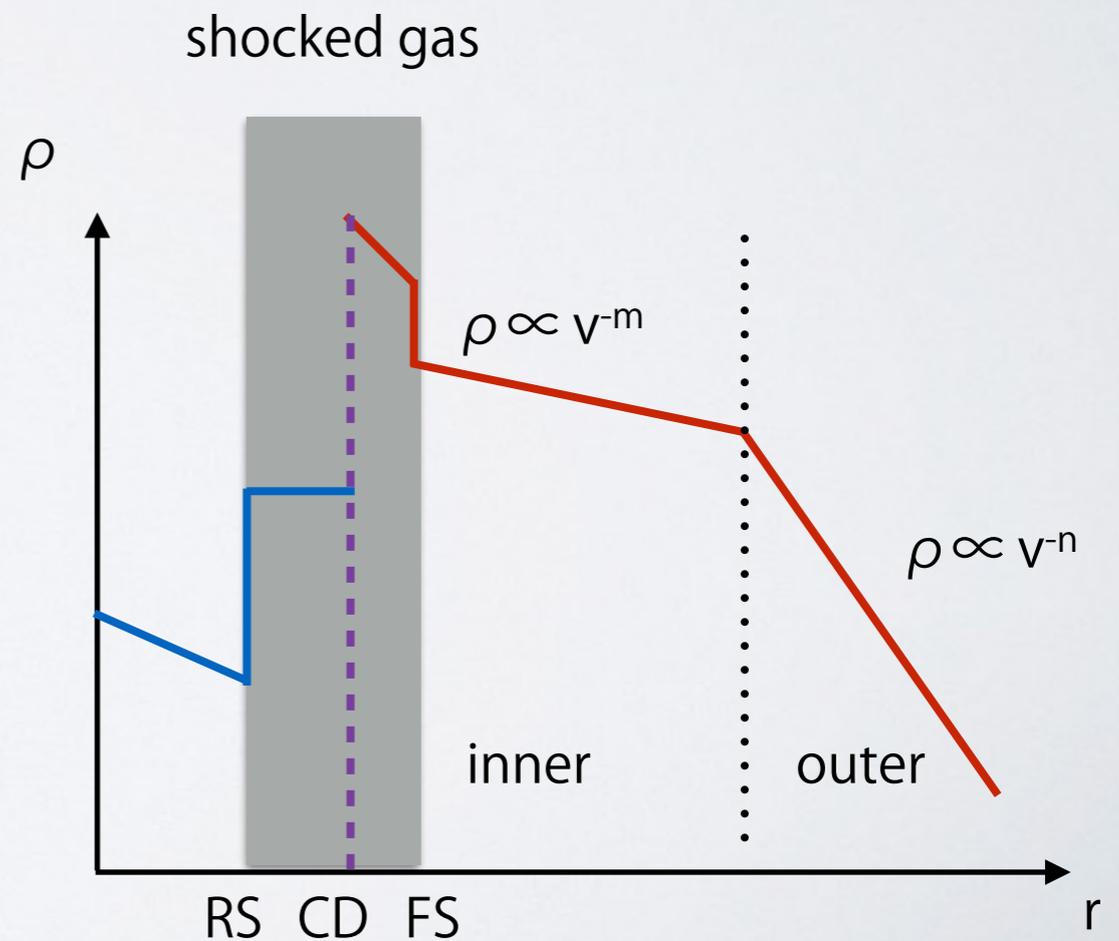
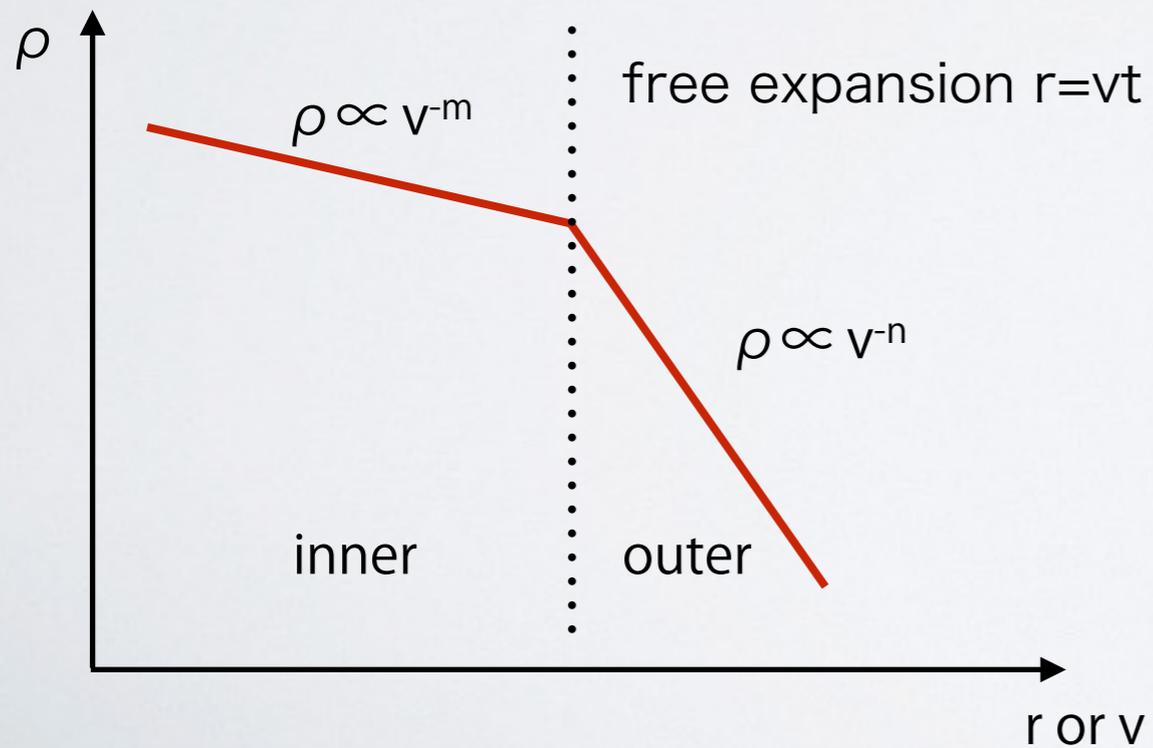
e.g., Chevalier (1992), Jun (1998)



freely expanding supernova ejecta (high ρ)

geometrically thin shell driven by hot bubble
assume $R_c \propto t^a$

high pressure region filled with hot gas (low ρ)
 $\rho \propto Lt/V_c$



1D analytic model

e.g., Chevalier (1992), Jun (1998)

pressure
of the hot bubble

$$p_c = \frac{3(\gamma - 1)E_{\text{th}}}{4\pi R_c^3} = \frac{3(\gamma - 1)(2 - \gamma)}{1 + 3\alpha(\gamma - 1)} \frac{Lt}{4\pi R_c^3}$$

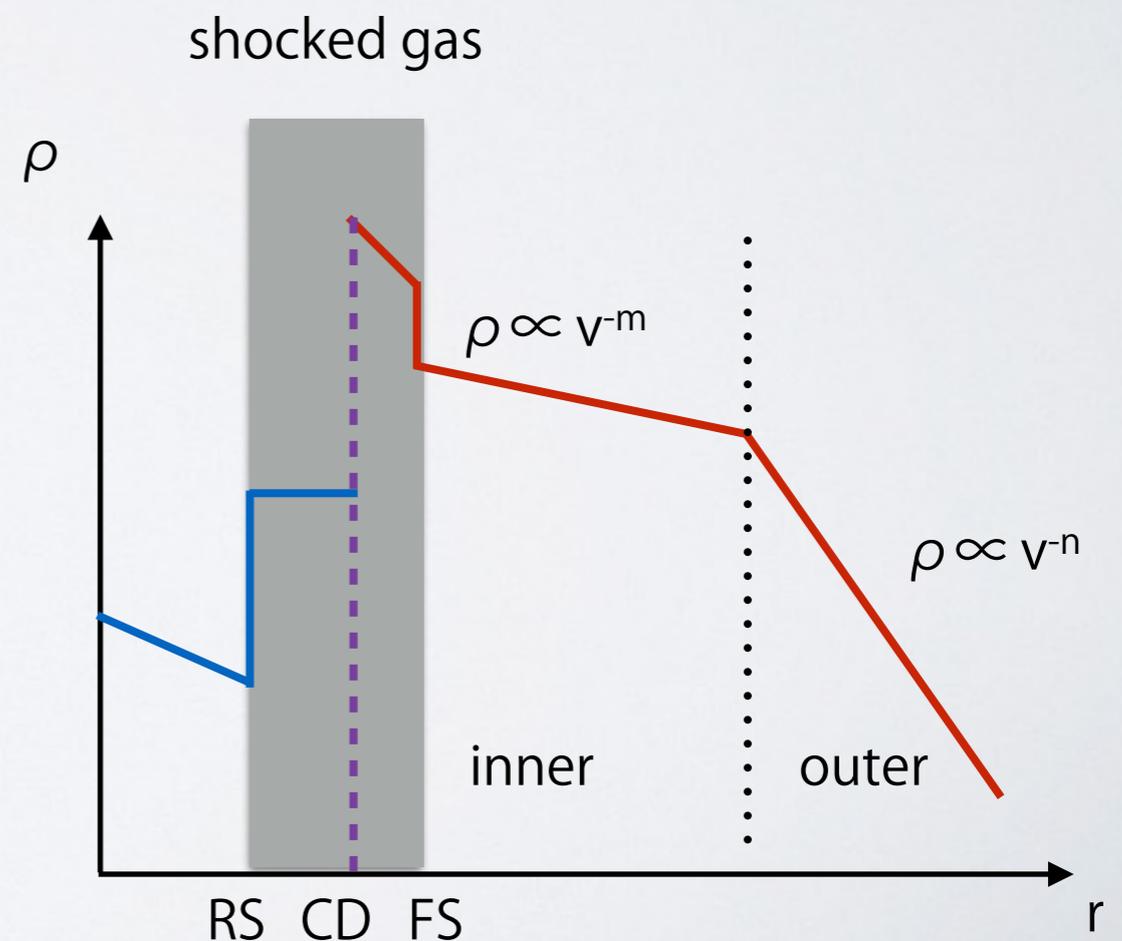
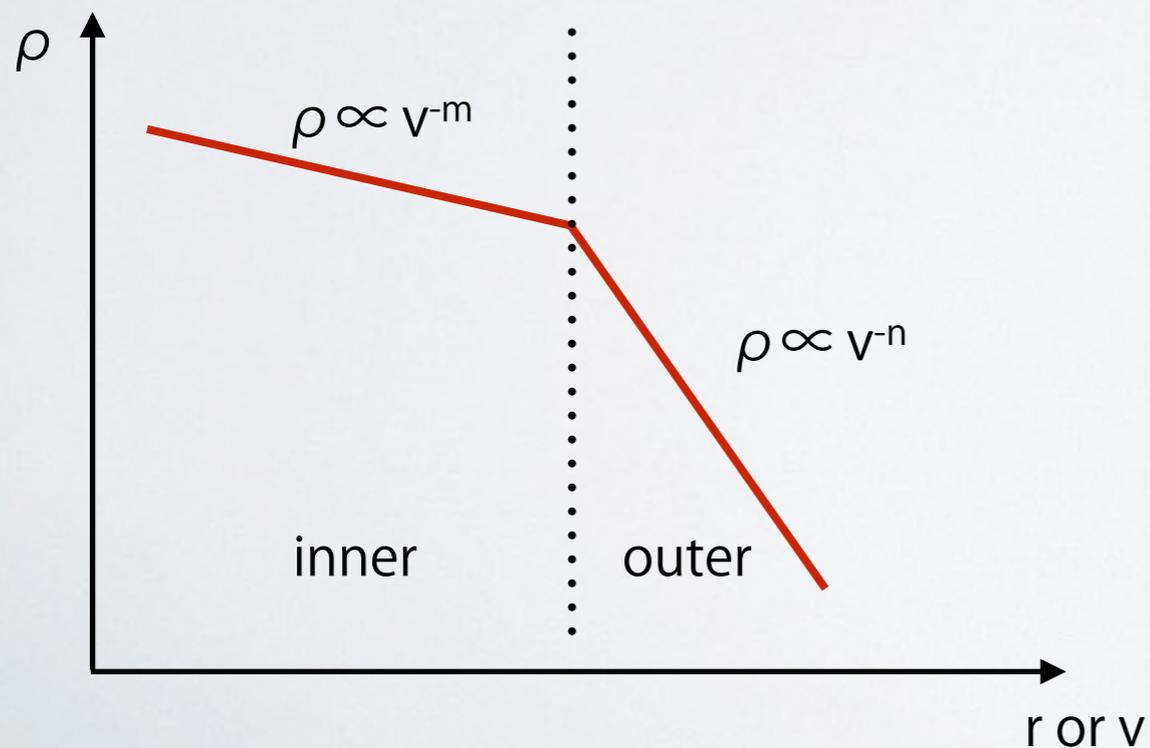
Eq. of continuity

$$M(t) \propto \int_0^r \rho(t, r)r^2 dr \propto t^{m-3}r^{3-m}$$

Eq. of motion

$$M(t) \frac{d^2 R_c}{dt^2} = 4\pi R_c^2 p_c \propto \frac{t}{R_c}$$

$$R_c^{5-m} \propto t^{6-m} \Rightarrow R_c \propto t^\alpha, \quad \text{with } \alpha = \frac{6-m}{5-m}$$



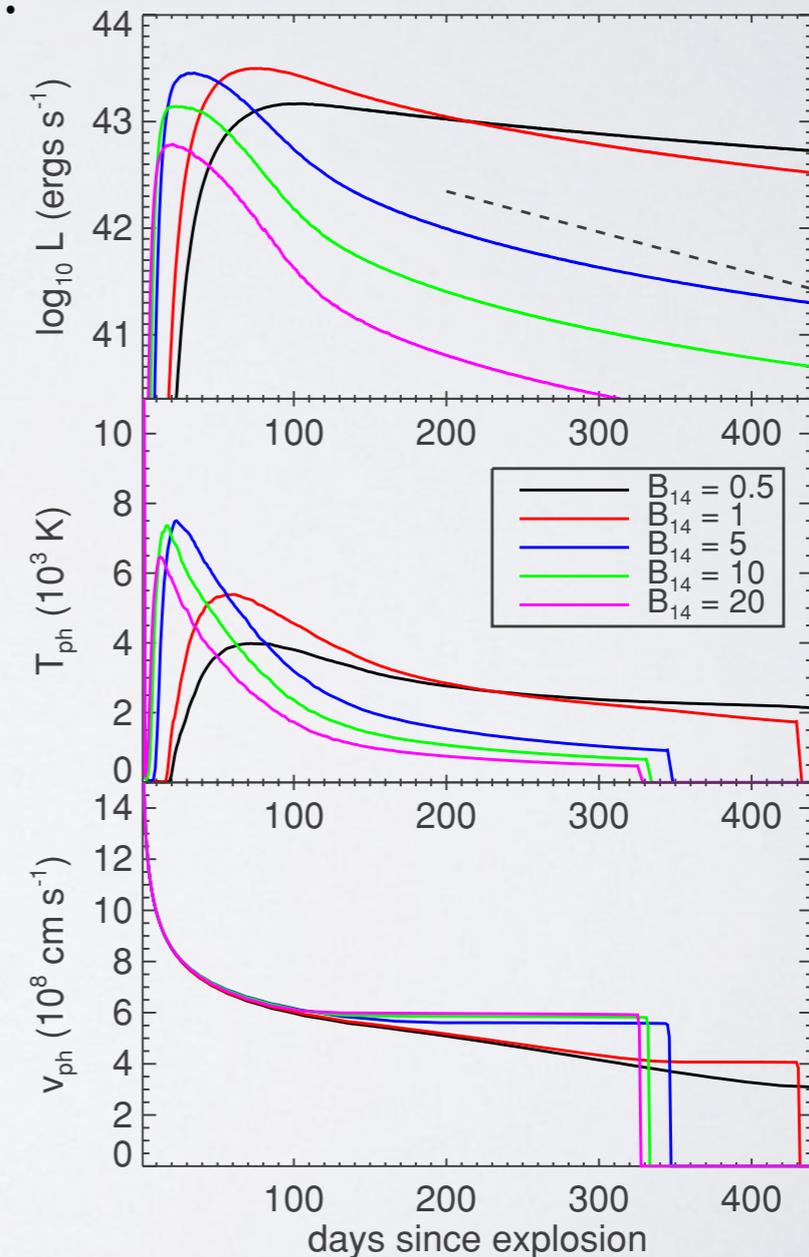
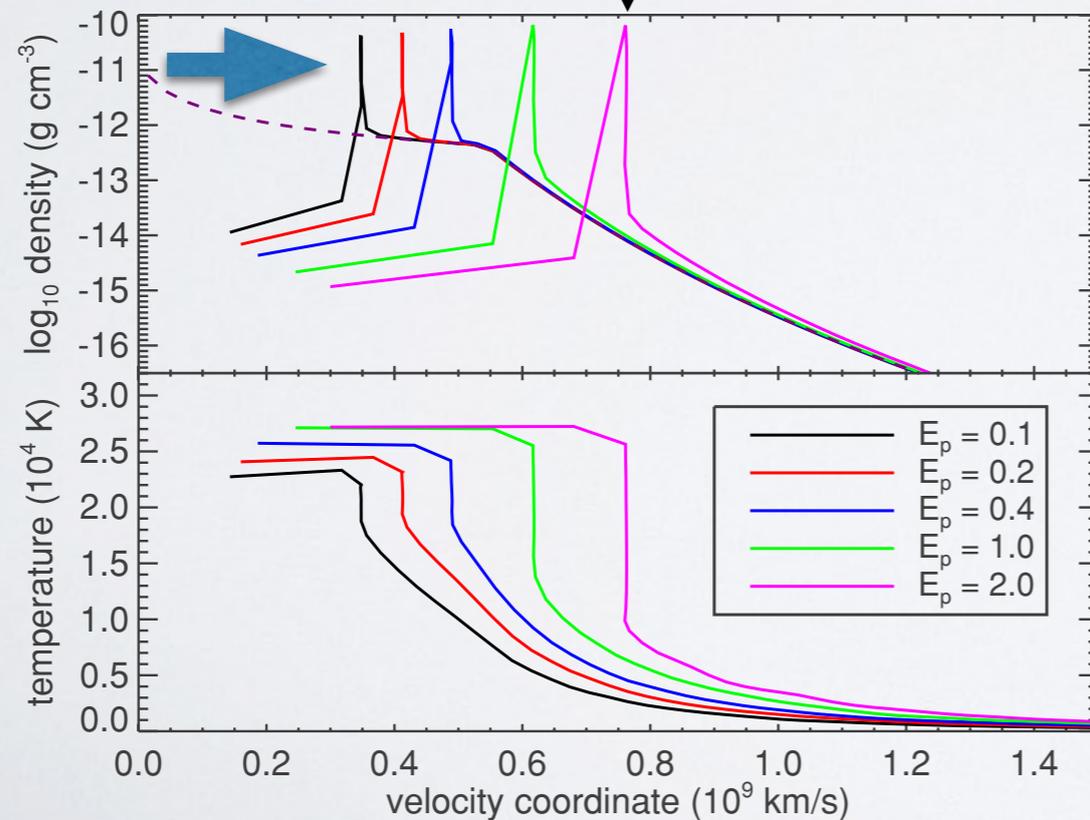
1D spherical picture of SN-wind interaction

- In 1D spherical case, the energy redistribution is realized by a geometrically thin shell (swept up materials by the energy injection) and the radiation diffusing out from the shell.
- It seems OK to explain the high brightness of SLSNe.

1D RHD simulation by Kasen&Bildsten (2010)

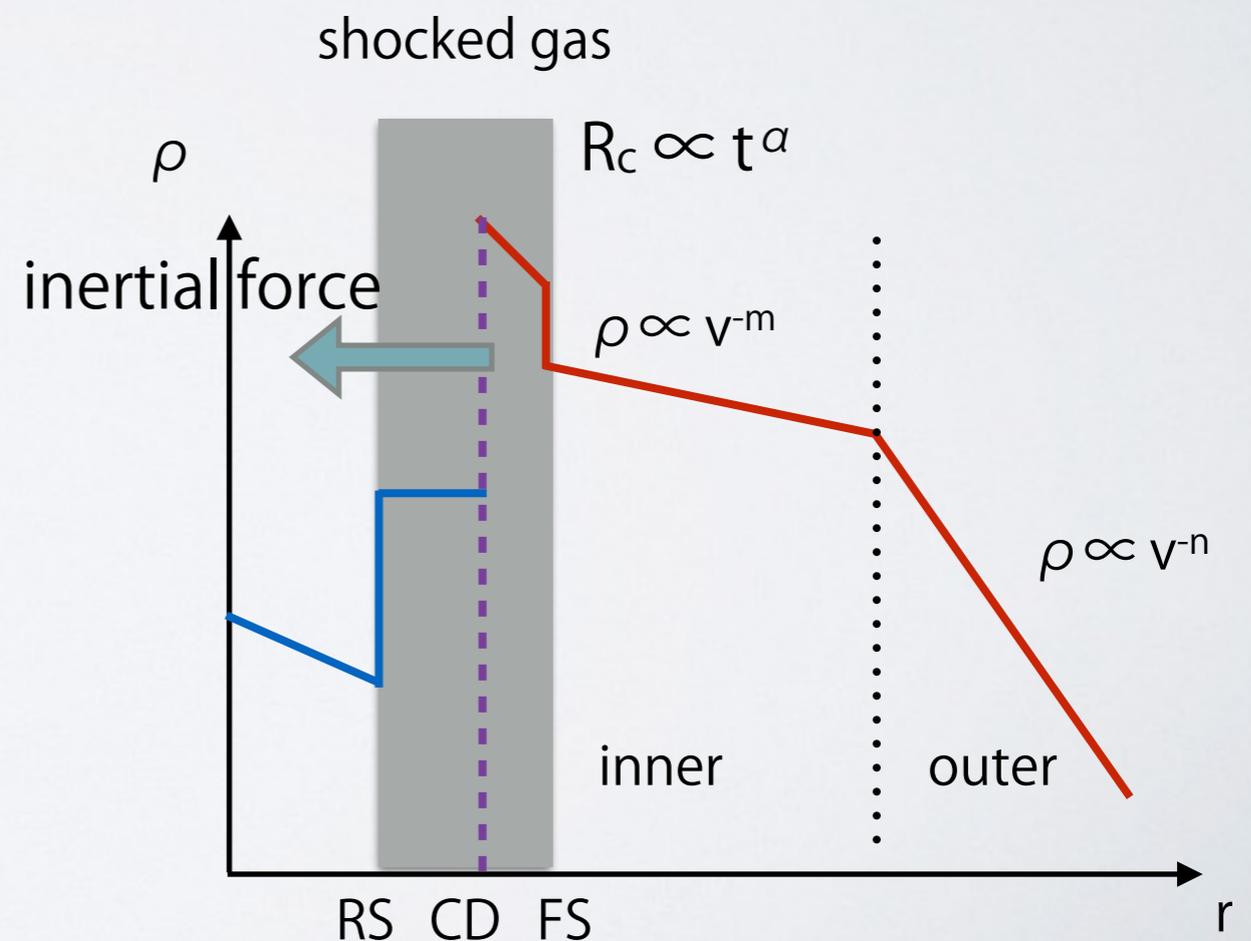
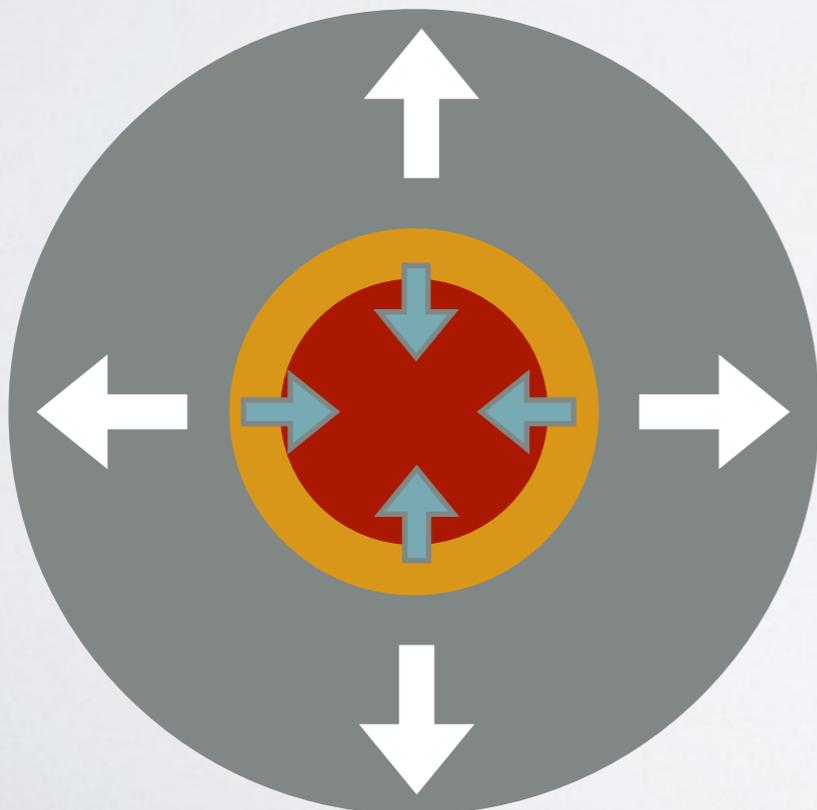
geometrically thin shell

magnetar energy injection



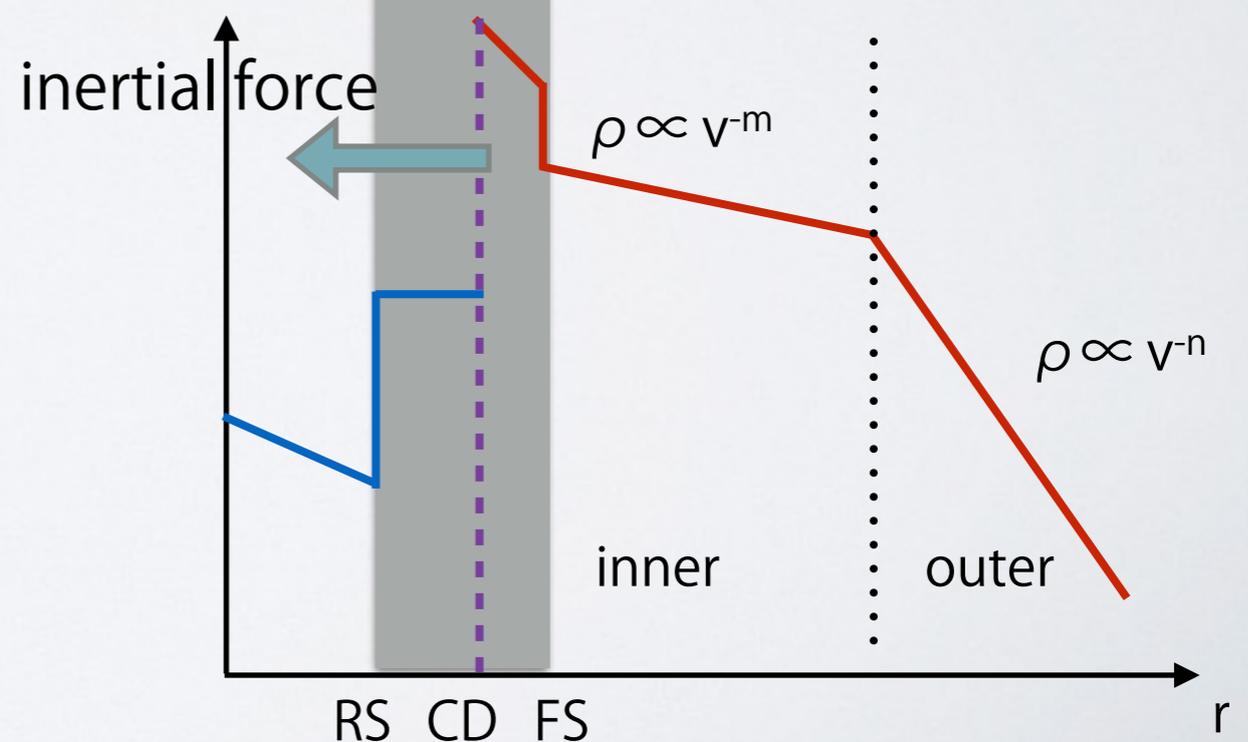
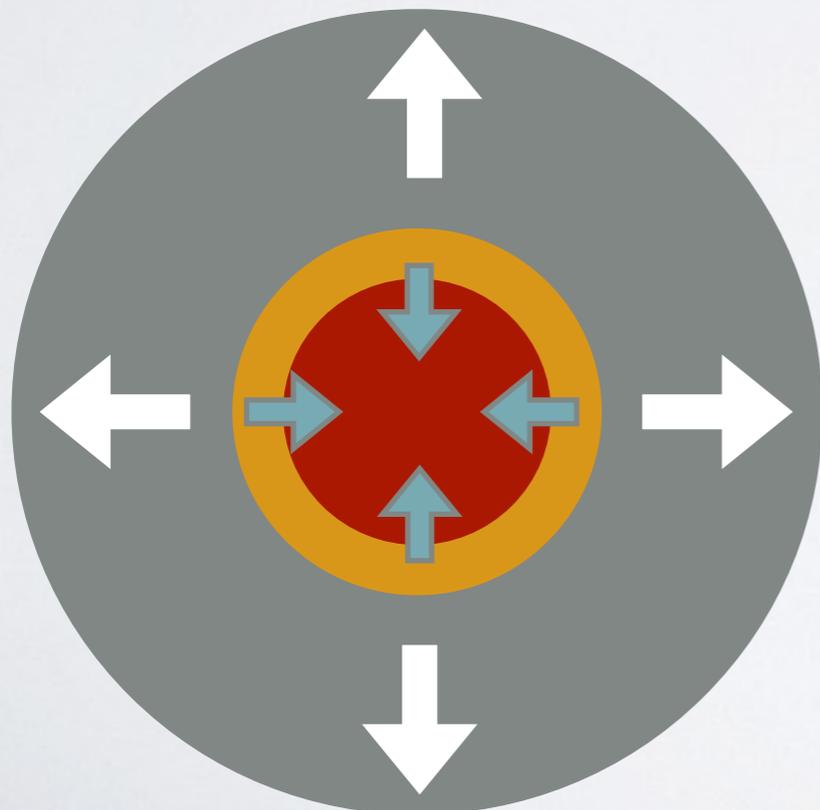
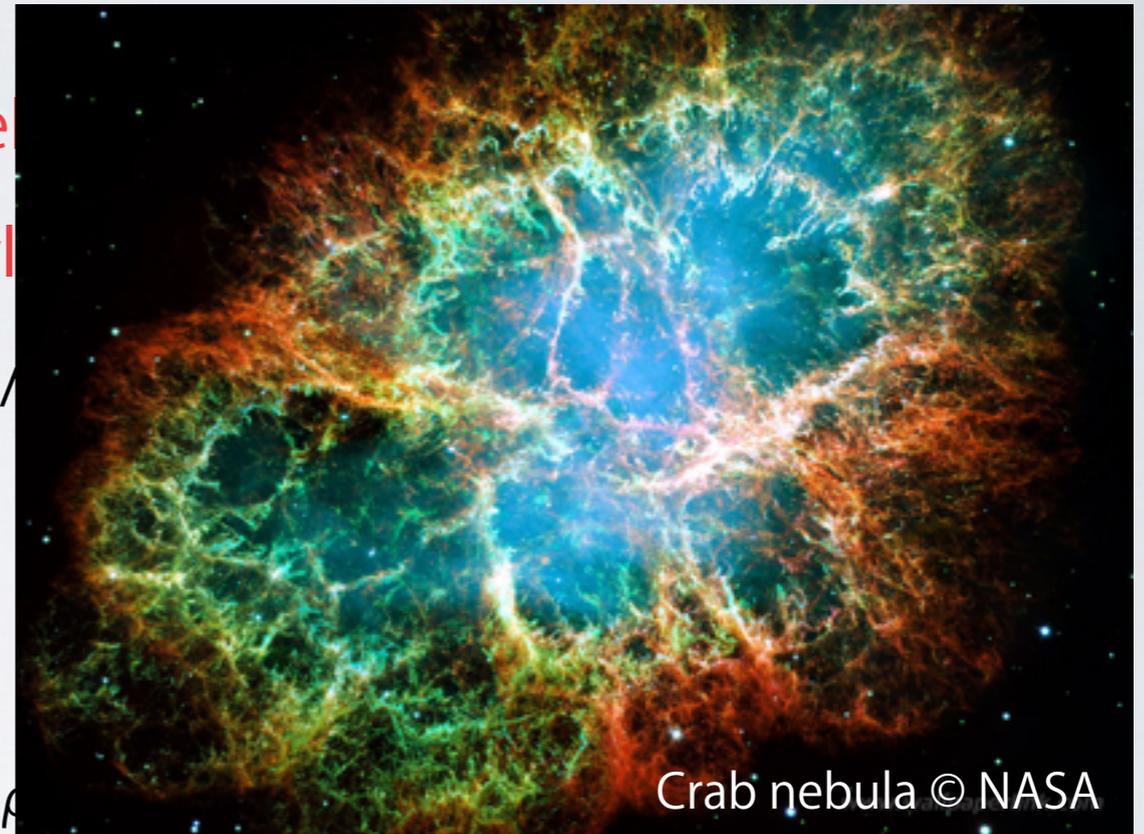
Q: Is 1D picture correct?

- No.
- From 1D analysis, we see that the shell is **accelerating ($\alpha > 1$)**
- an accelerating spherical shell is **Rayleigh-Taylor unstable!**
- more precisely, the unstable condition is “ $(dp/dr) \times (d\rho/dr) < 0$ ”



Q: Is 1D picture correct?

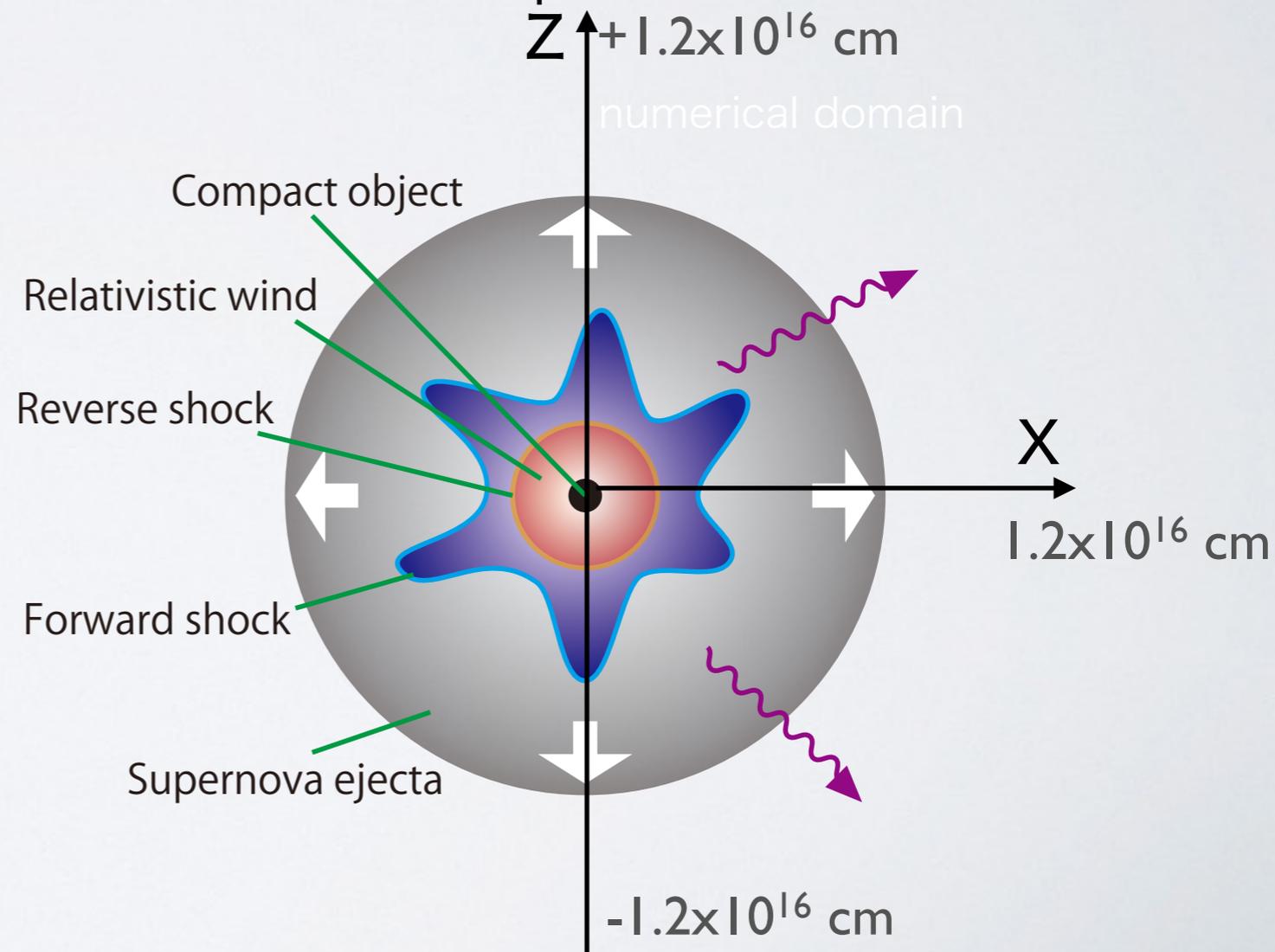
- No.
- From 1D analysis, we see that the shell is **accelerating**
- an accelerating spherical shell is **Rayleigh-Taylor unstable**
- more precisely, the unstable condition is “ $(dp/dr) < \rho a$ ”



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

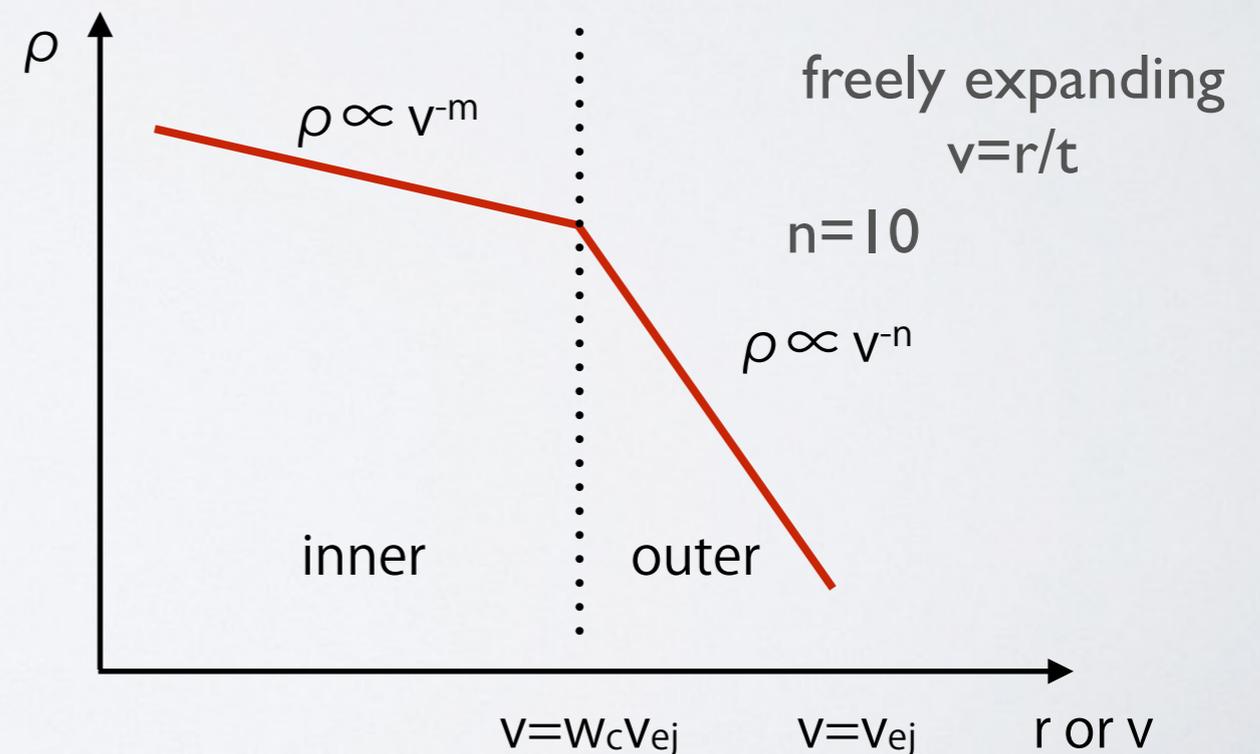
- cartesian coordinate (x,y,z)
- x,y,z in $[0, 1.2 \times 10^{16} \text{ cm}]$
- AMR technique.
- ideal gas law $\gamma = 4/3$
- relativistic gas injection within $3 \times 10^{12} \text{ [cm]}$: $L = 10^{46} \text{ [erg/s]}$ up to 10^{52} [erg]
- $dM/dt = 0.05L/c^2$
- SN ejecta with $10[M_{\odot}]$ and 10^{51} [erg]
- unit time $t_c = E_{sn}/L = 10^5 \text{ sec}$
- from $t = 0.1t_c$ up to $t = 20.0t_c$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

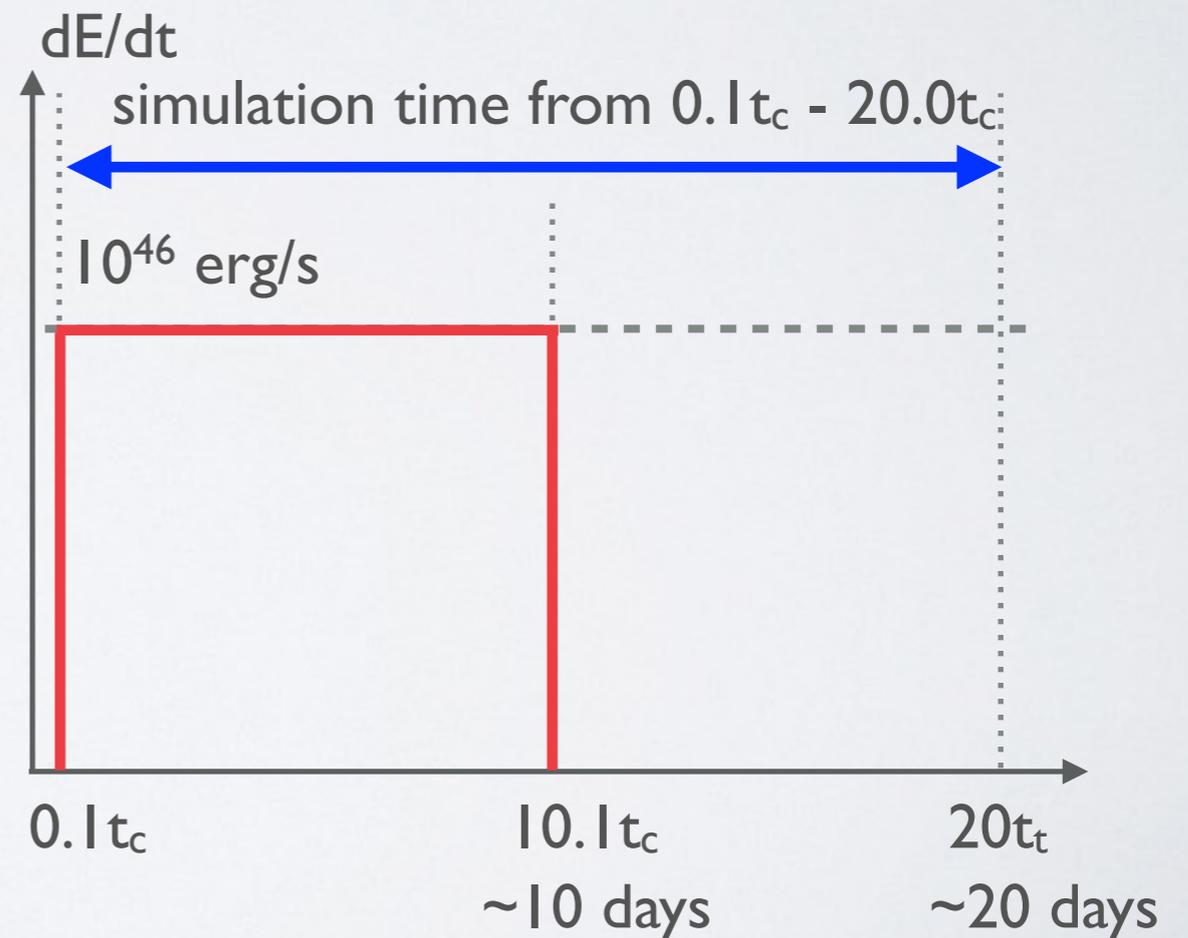
- cartesian coordinate (x,y,z)
- x,y,z in $[0, 1.2 \times 10^{16} \text{ cm}]$
- AMR technique.
- ideal gas law $\gamma = 4/3$
- relativistic gas injection within $3 \times 10^{12} \text{ [cm]}$: $L = 10^{46} \text{ [erg/s]}$ up to 10^{52} [erg]
- $dM/dt = 0.05L/c^2$
- SN ejecta with $10[M_{\odot}]$ and 10^{51} [erg]
- unit time $t_c = E_{sn}/L = 10^5 \text{ sec}$
- from $t = 0.1t_c$ up to $t = 20.0t_c$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

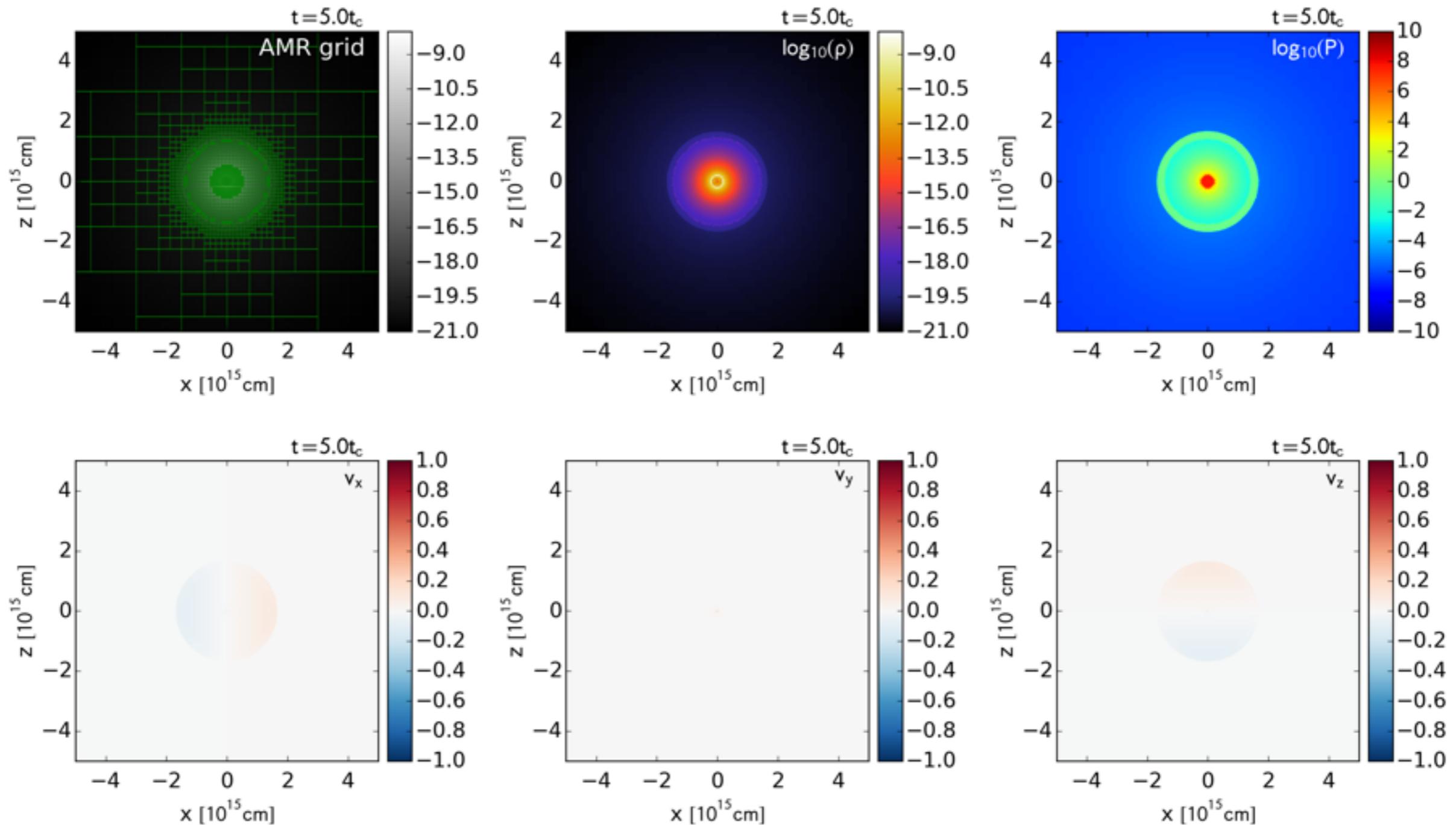
- cartesian coordinate (x,y,z)
- x,y,z in $[0, 1.2 \times 10^{16} \text{ cm}]$
- AMR technique.
- ideal gas law $\gamma = 4/3$
- relativistic gas injection within $3 \times 10^{12} \text{ [cm]}$: $L = 10^{46} \text{ [erg/s]}$ up to 10^{52} [erg]
- $dM/dt = 0.05L/c^2$
- SN ejecta with $10[M_{\odot}]$ and 10^{51} [erg]
- unit time $t_c = E_{\text{sn}}/L = 10^5 \text{ sec}$
- from $t = 0.1t_c$ up to $t = 20.0t_c$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

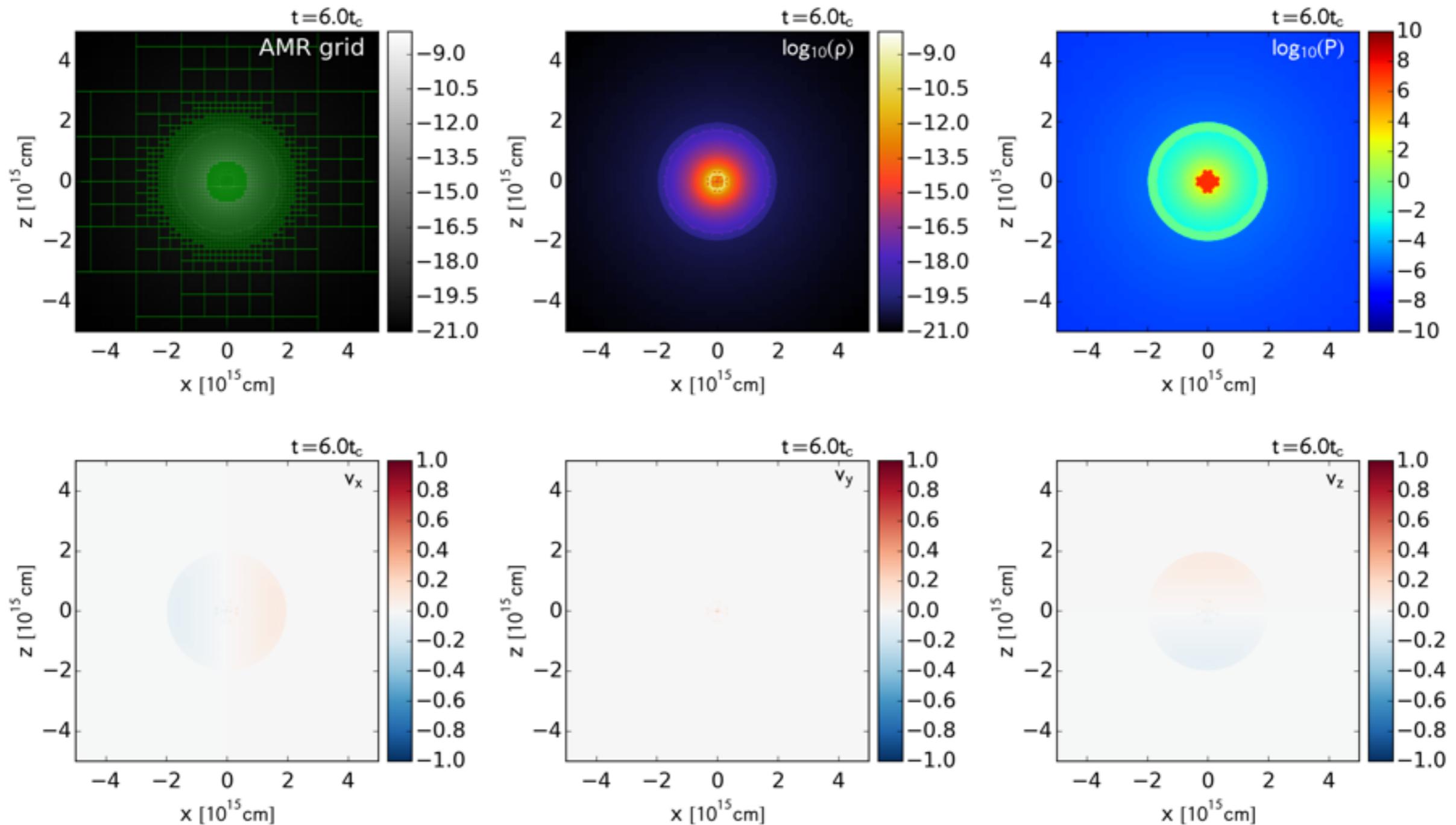
$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

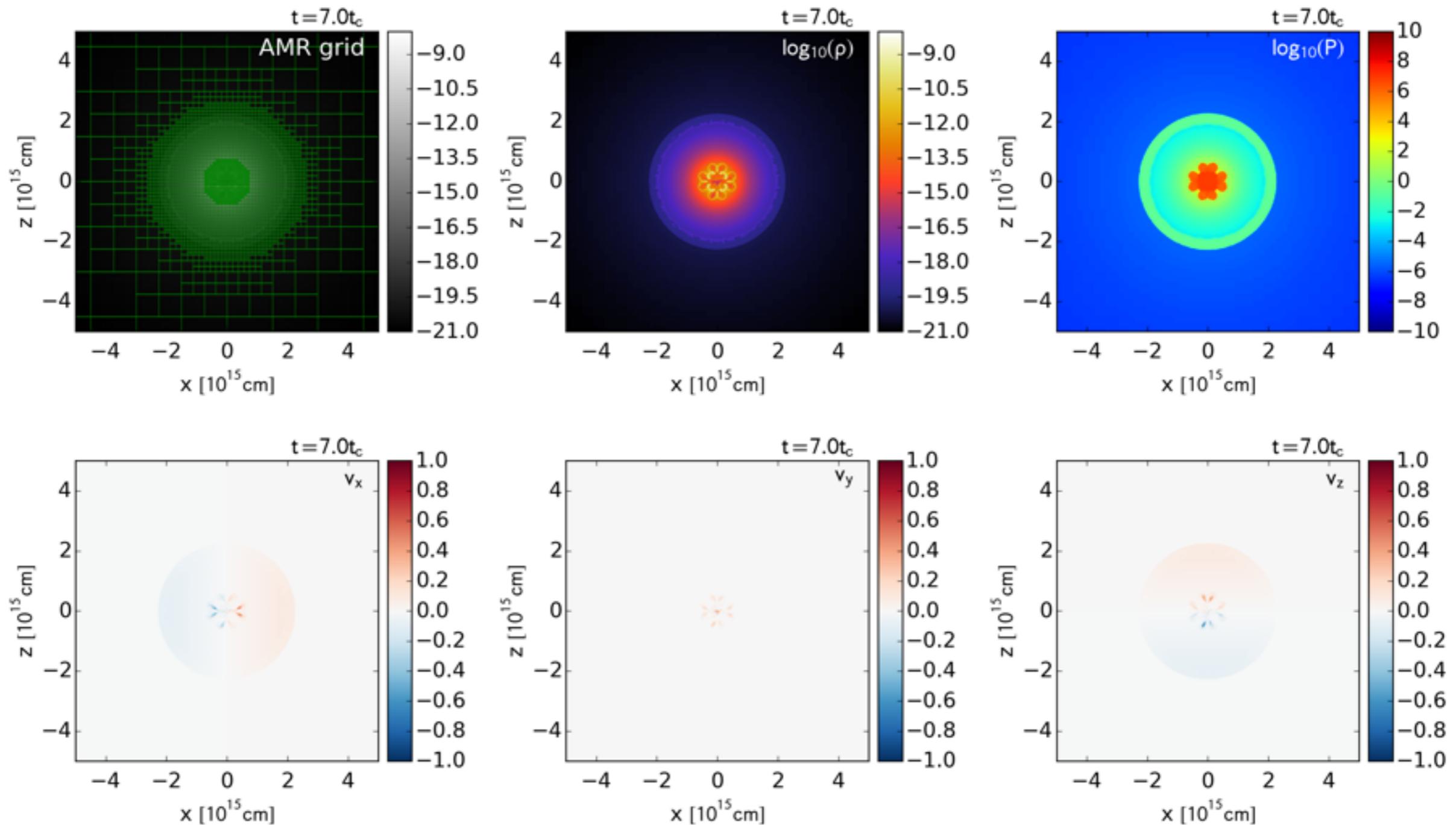
$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

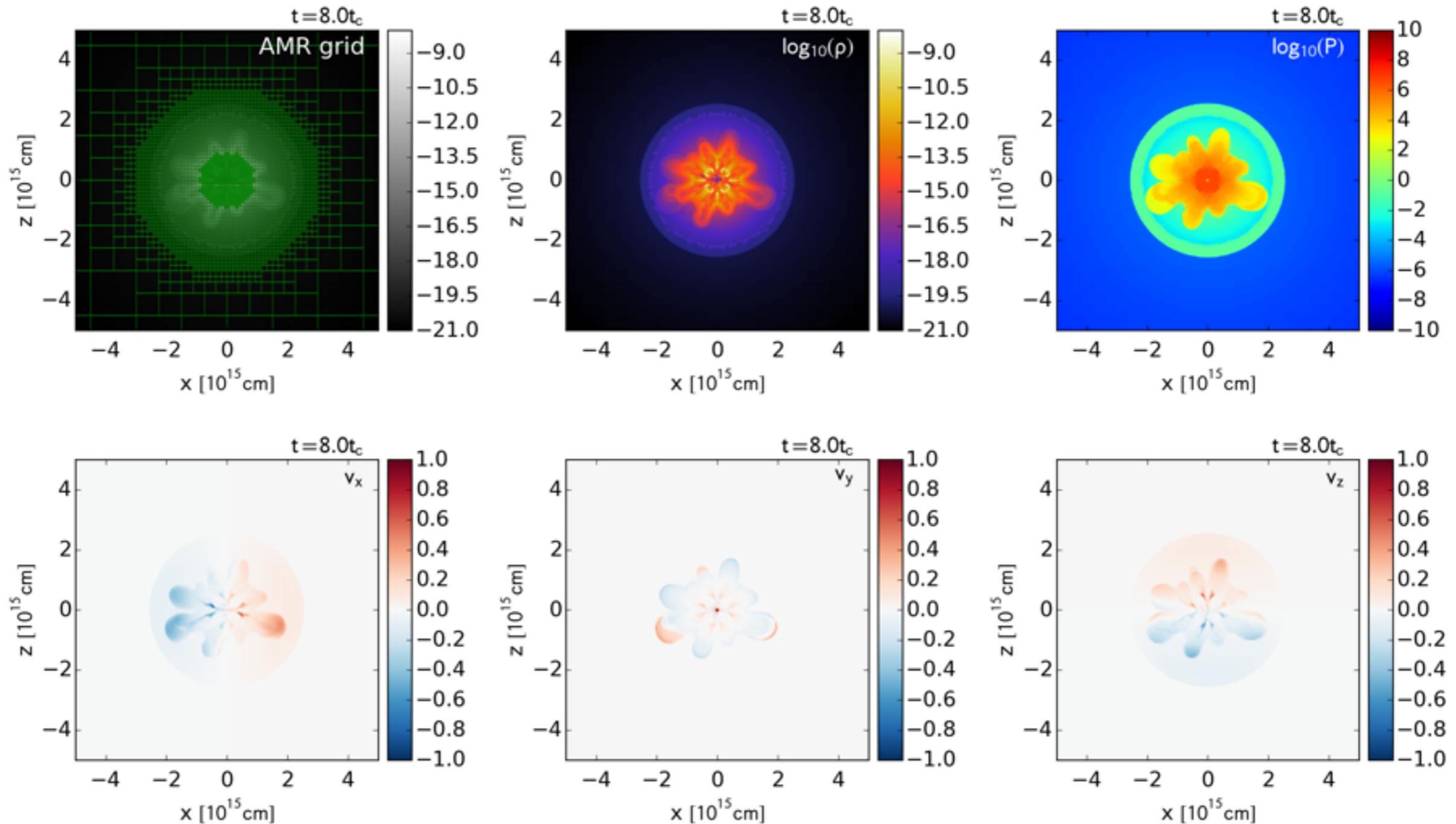
$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

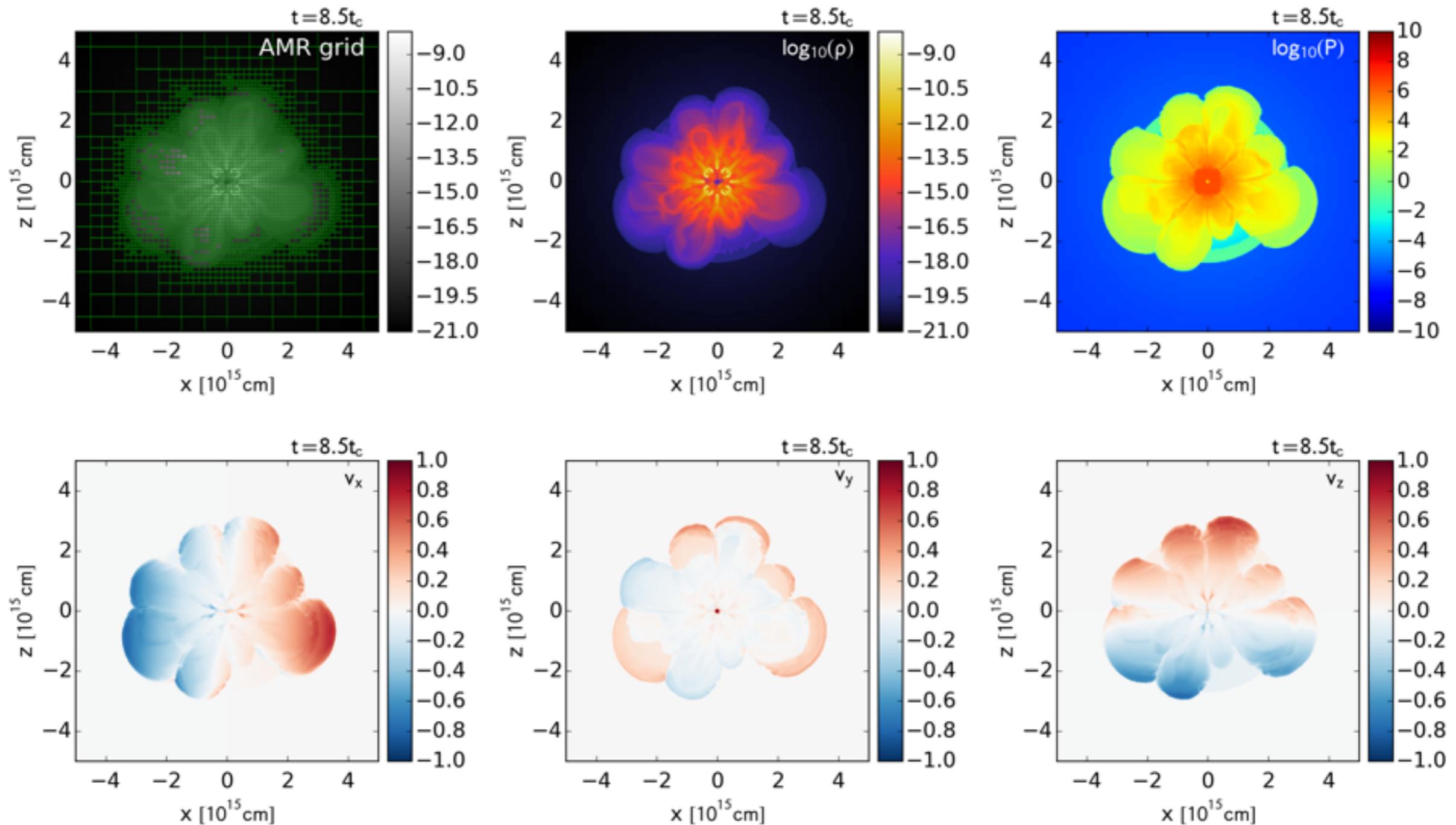
$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

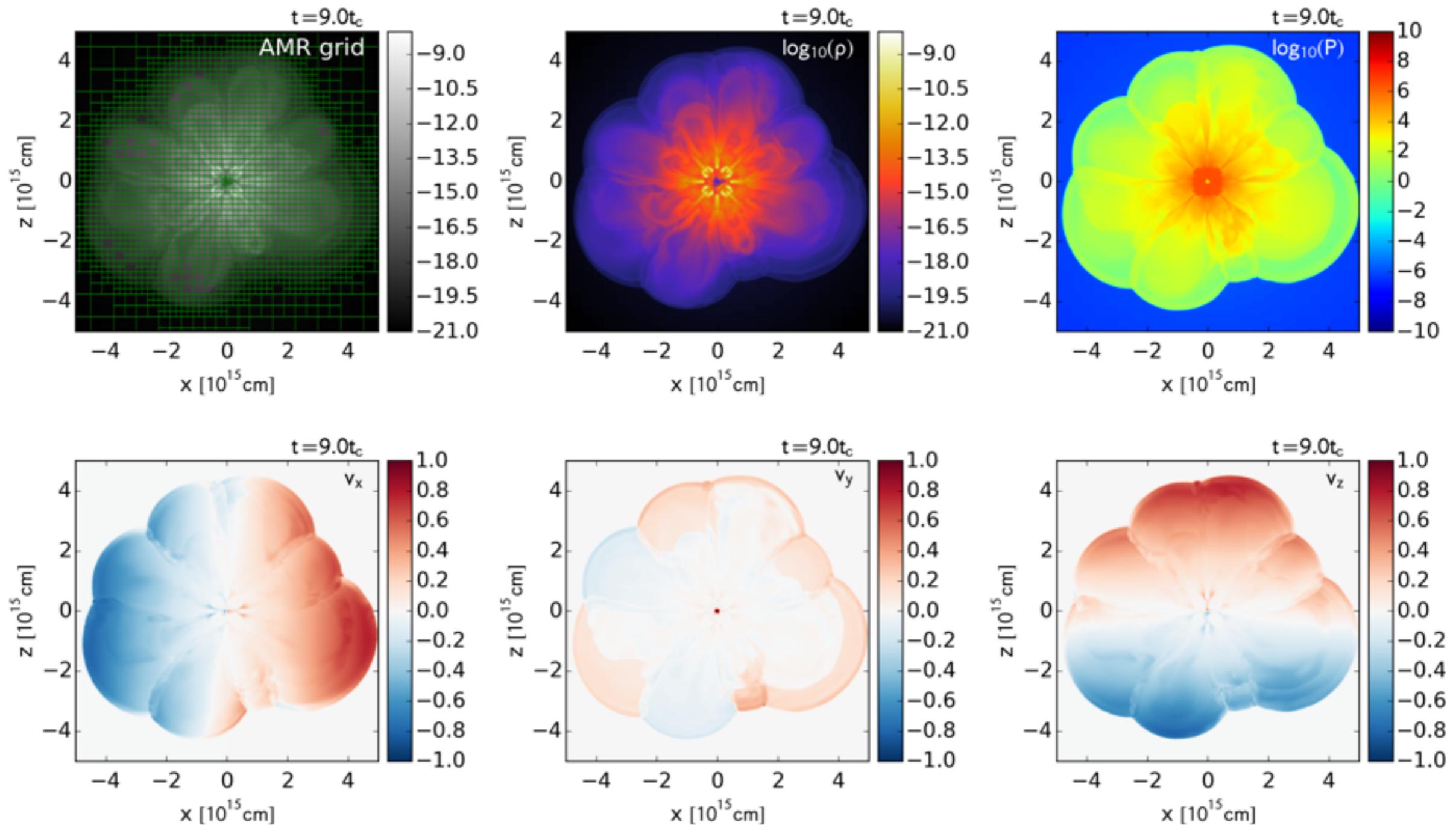
$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

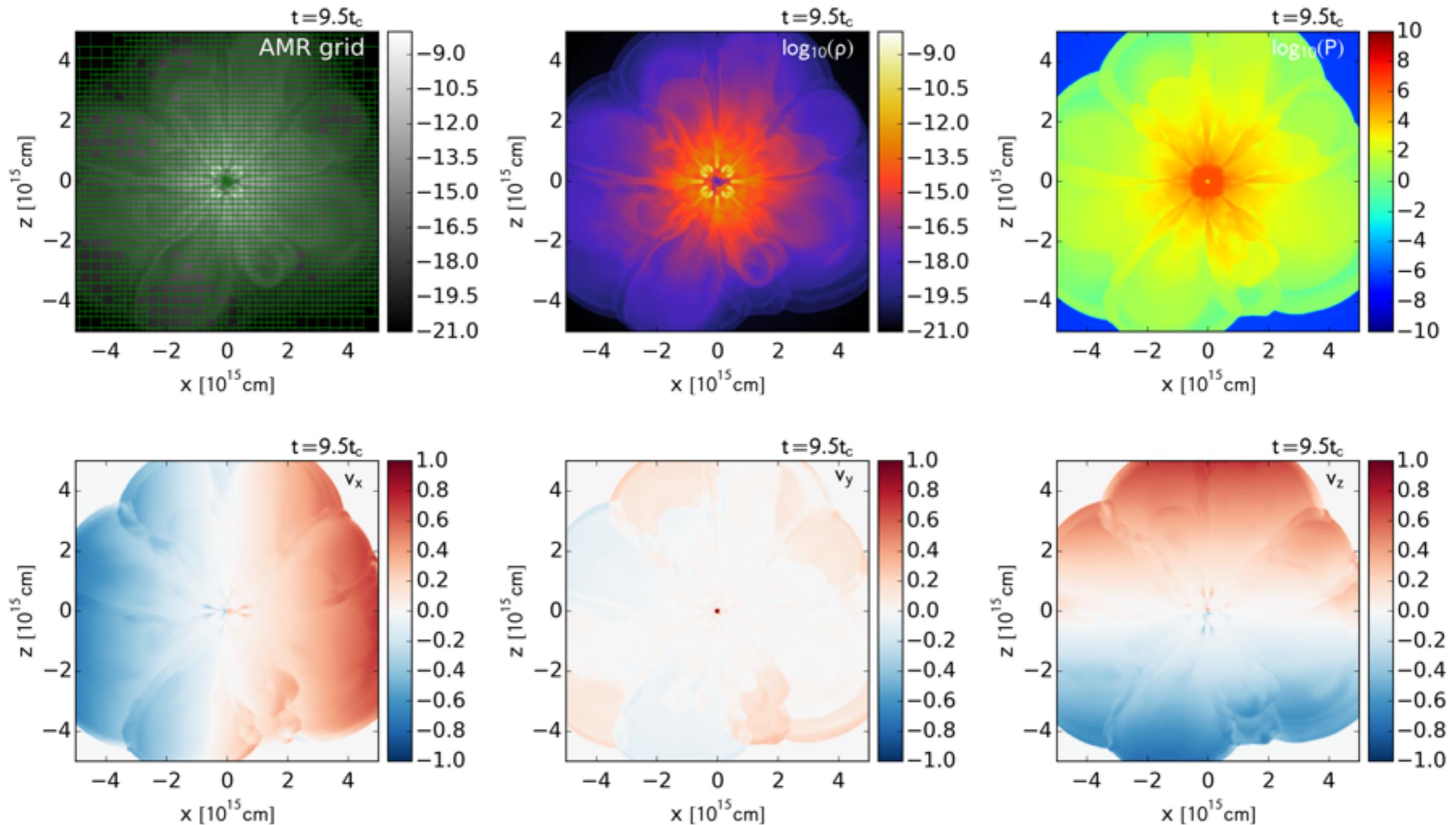
$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$



3D simulation (Suzuki&Maeda, in prep.)

cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

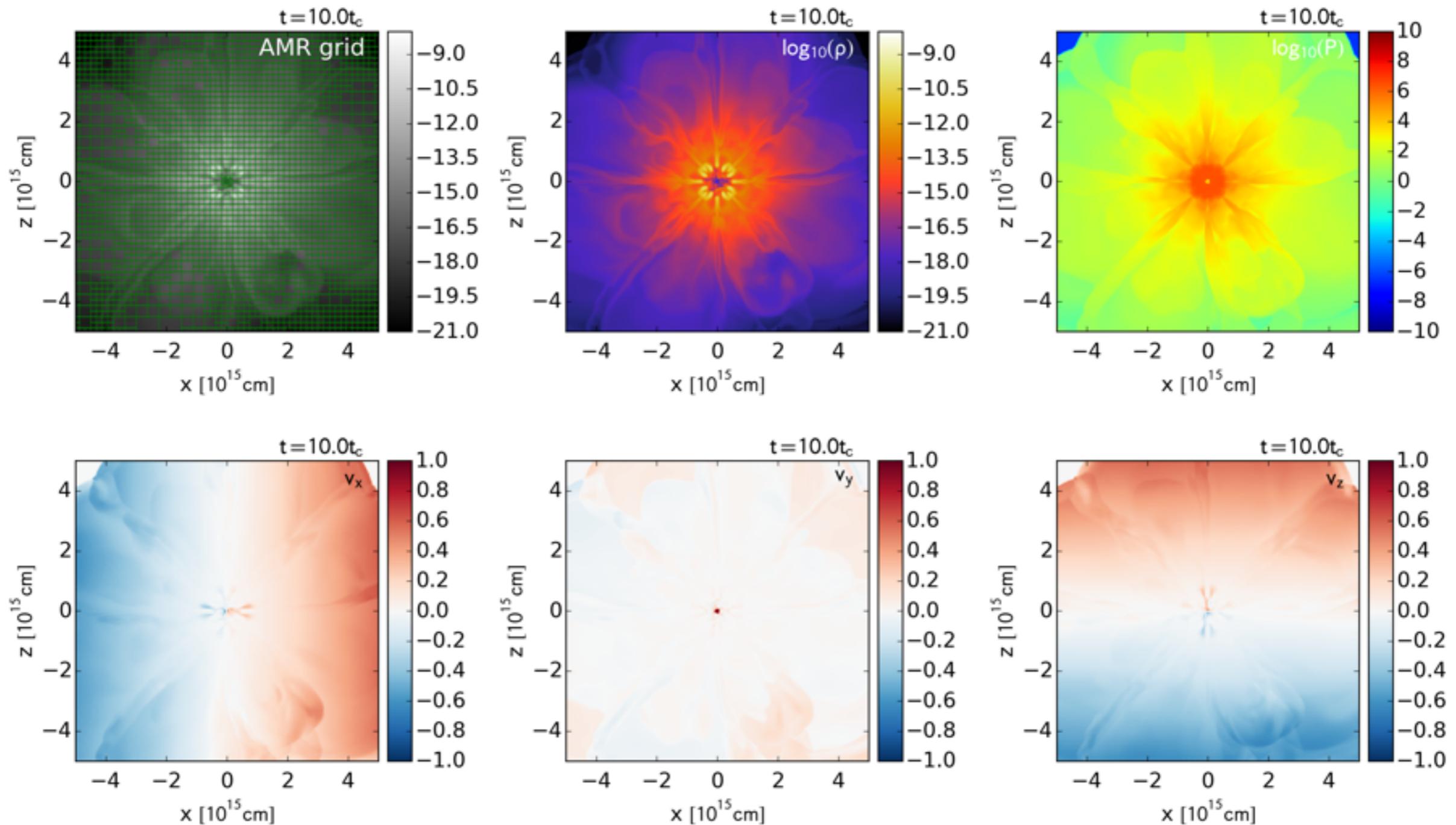
$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$



3D simulation (Suzuki&Maeda, in prep.)

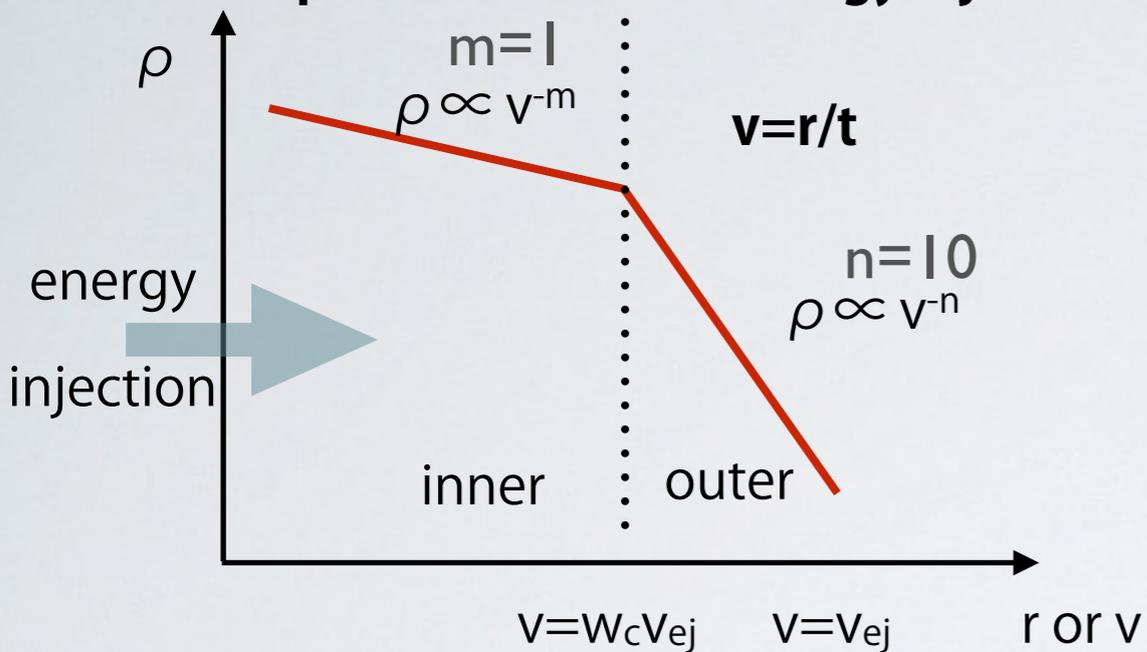
cf. 2D cylindrical simulation (Suzuki&Maeda 2017)

$$E_{\text{sn}}=10^{51} \text{ [erg]}, \quad L=10^{46} \text{ [erg/s]}, \quad t_c=10^5 \text{ [sec]}$$

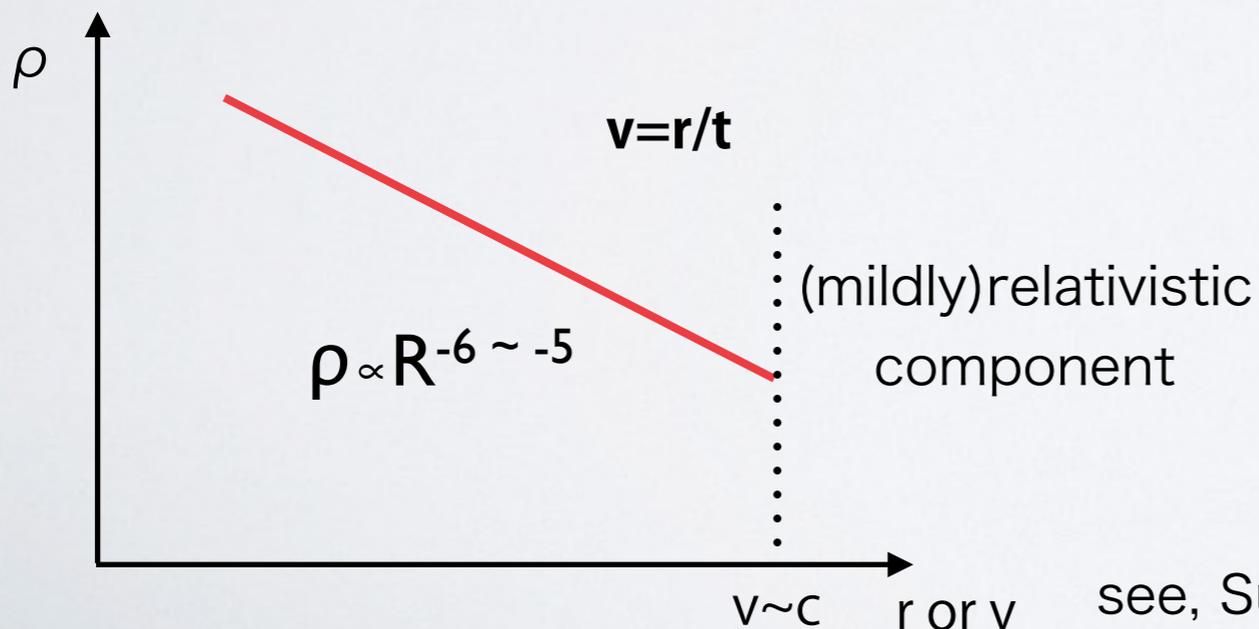


Density structure

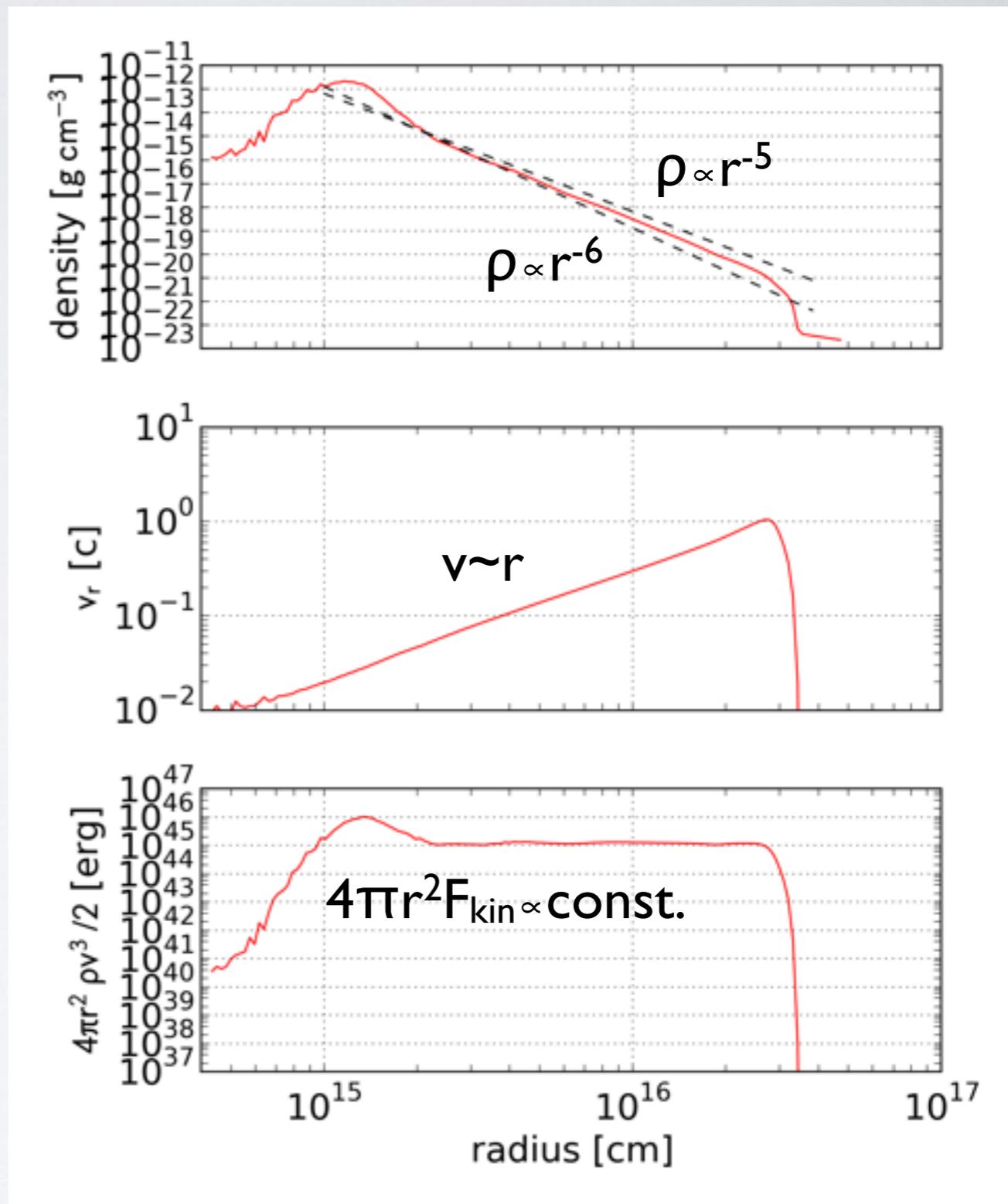
Free expansion "before" energy injection



Free expansion "after" energy injection



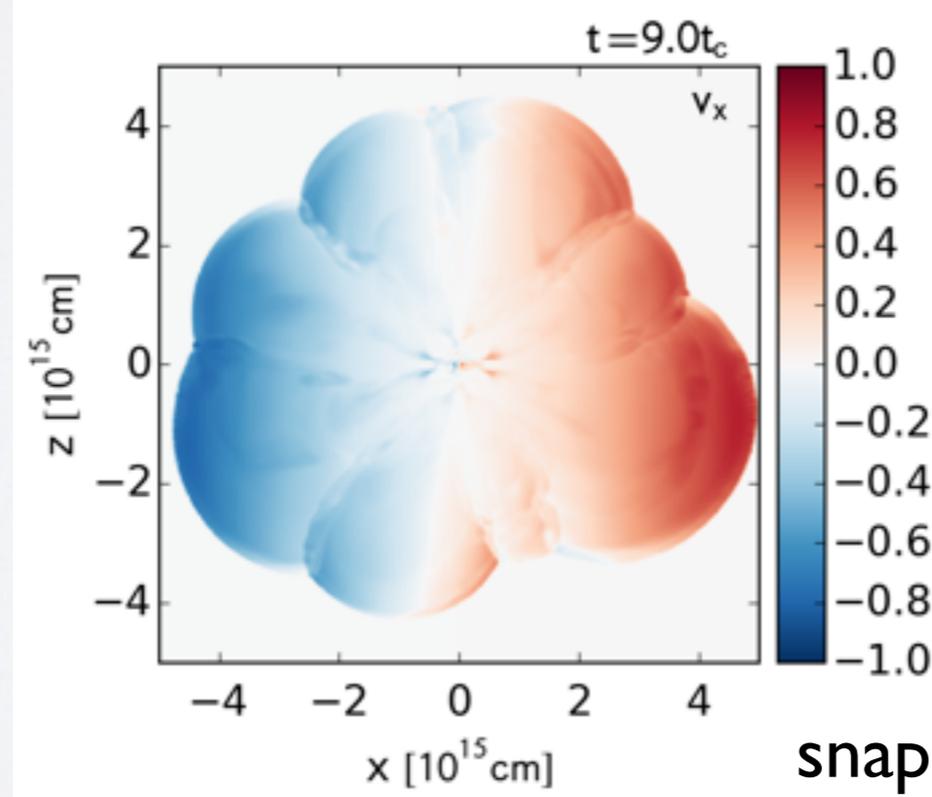
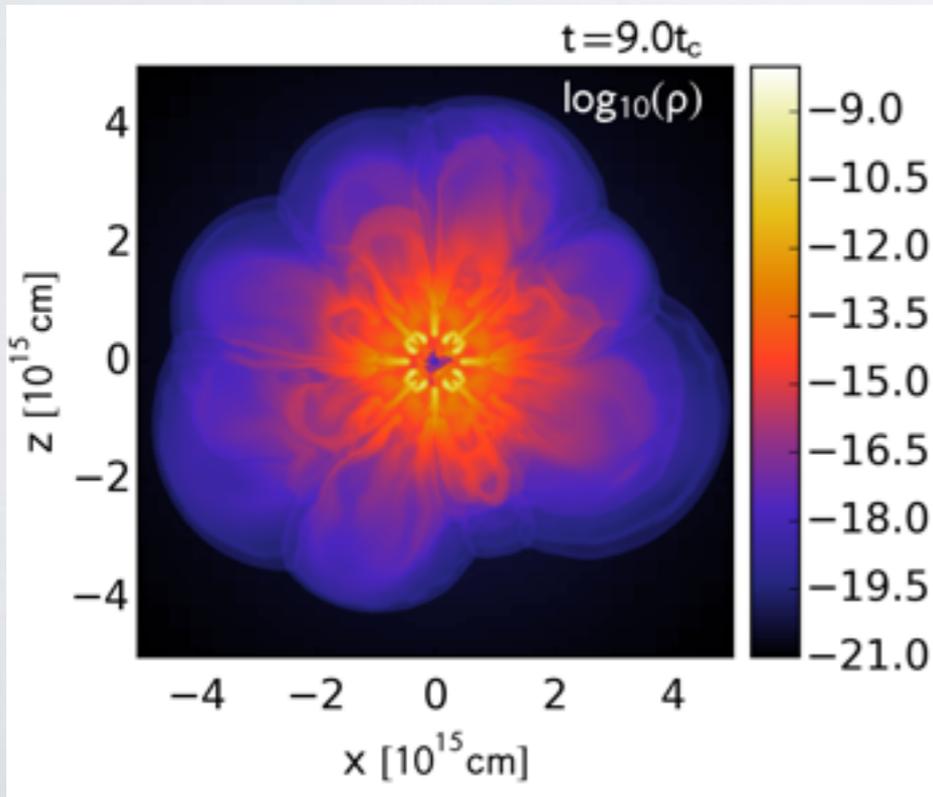
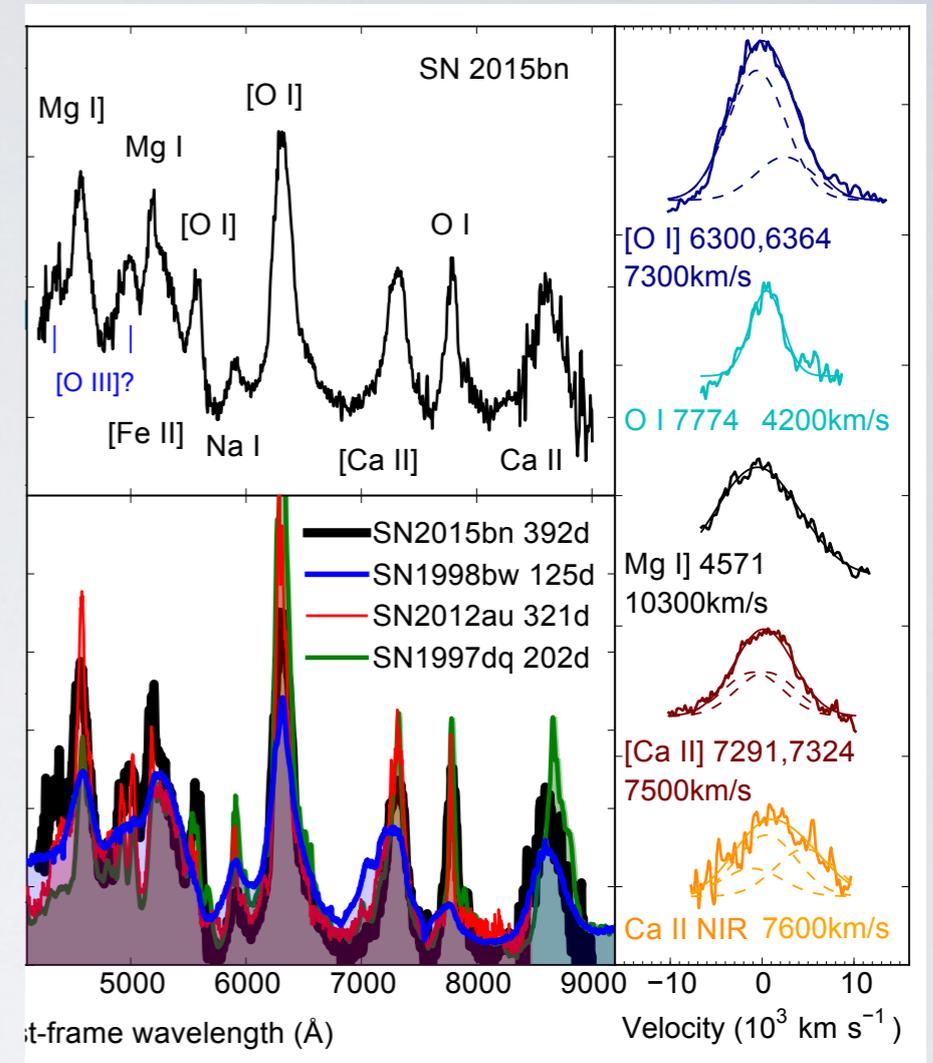
radial profiles at $t=20t_c$



see, Suzuki&Maeda (2017)

Density structure

- shredding by hot bubble breakout
- efficient matter mixing + high photospheric velocity
- broad-line feature in optical spectrum(?)
- $\rho \propto v^{-7}$ is favored for SLSNe-I (Mazzali+2016)



Nicholl+ (2016)

↑
?

今後、可視光
スペクトルを計算

snapshot at $t=9t_c \sim 9$ days

Summary: central-engine SNe in multi-D

- Dynamical evolution of SN ejecta + additional energy injection is multi-dimensional
- Hot bubble breakout leads to violent mixing (+ precursor?)
- final radial density structure of the ejecta is a simple power-law function
- we have started 3D simulations and confirmed the picture

